

# Axion Results with & Future Searches with



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on behalf of the LUX and LZ Collaborations



*Institute of Physics Joint APP and HEPP Annual Conference, University of Sheffield, 11/04/2017*

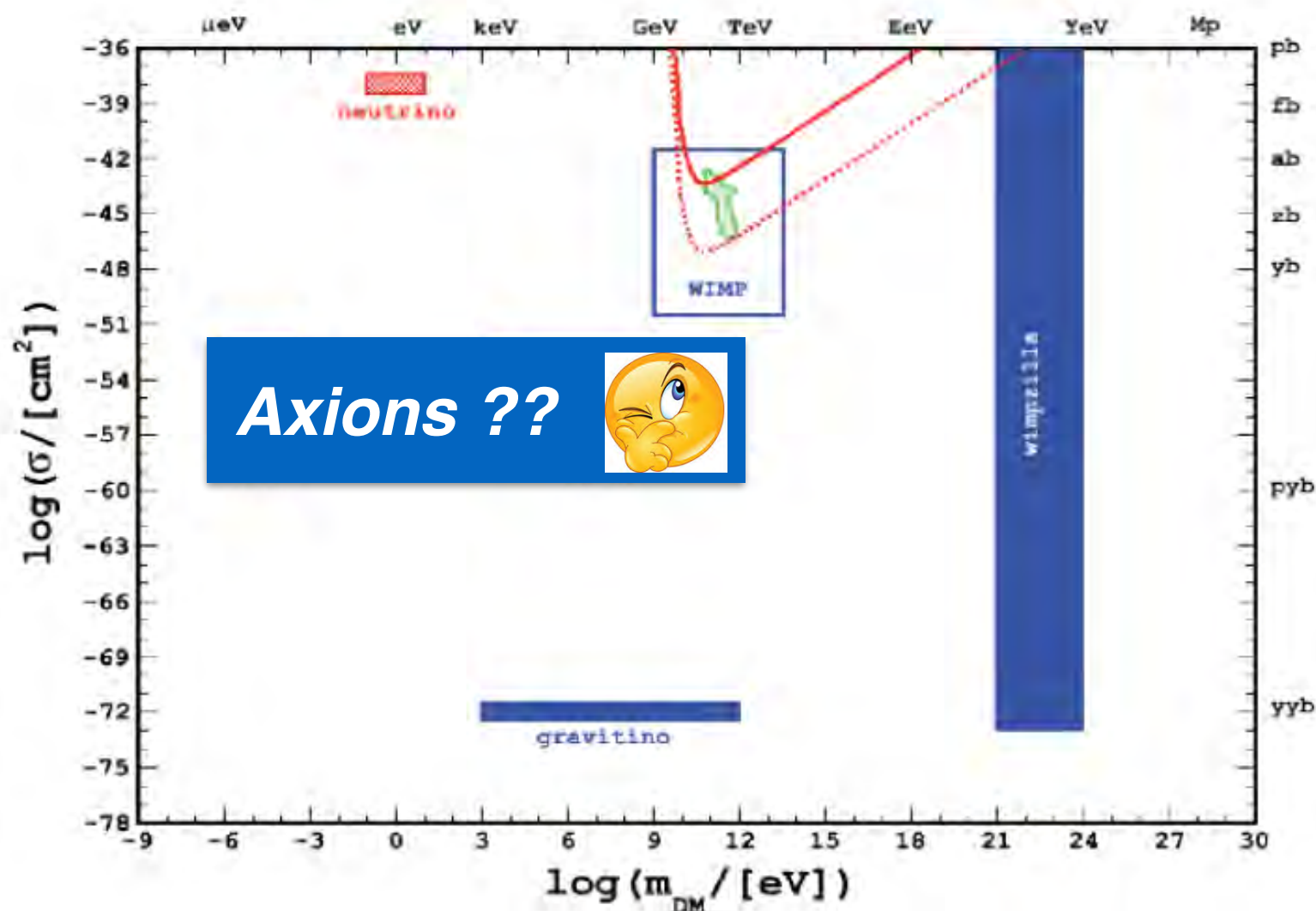
# Outline

- Axions: why, where, how
- Axion searches with xenon time projection chambers
- First axion results with LUX
- Sensitivity for axion searches with LUX-ZEPLIN
- Summary



# Why axions ?

- In Particle Physics, the axion field provides a dynamical solution to the strong CP violation problem (Peccei-Quinn solution)
- Axions do have the main DM characteristics: nearly collisionless, neutral, non baryonic, present within the Universe in sufficient quantities to provide the DM density

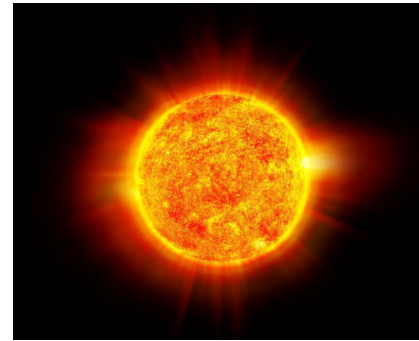


- Extensions of the Standard Model of Particle Physics introduce the so called axion-like particles (ALPs), which could be dark matter candidates
- The scenario of Dark Matter searches can be wider than just WIMPs

# Where do axions come from?

- Potential sources of axions:

- **axions come from the Sun**

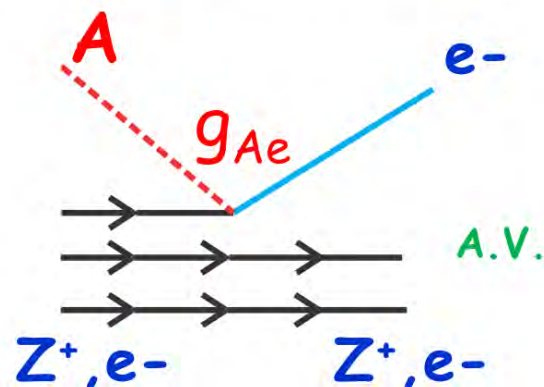


- **ALPs slowly move within our Galaxy**



## How do axions/ALPs couple?

Axio-electric effect



- Axions and ALPs can couple with electrons, via the so called **axio-electric effect**

- we can measure the coupling between axions/ALPs and electrons ( $g_{Ae}$ )

$$\sigma_{Ae} = \sigma_{pe}(E_A) \frac{g_{Ae}^2}{\beta_A} \frac{3E_A^2}{16\pi\alpha_{em}m_e^2} \left(1 - \frac{\beta_A^{2/3}}{3}\right)$$

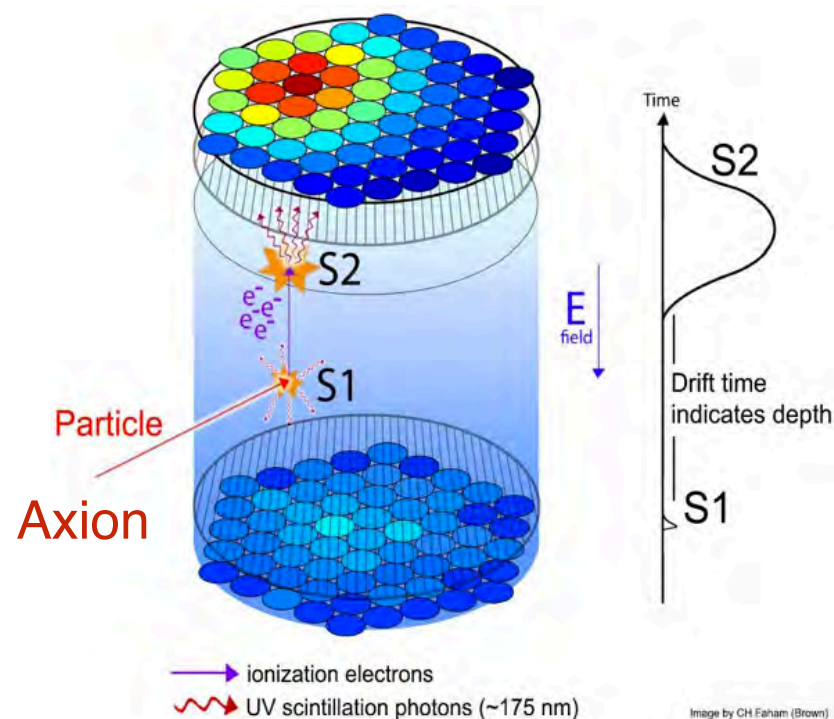
F. T. Avignone et al., Phys. Rev. D  
35, 2752 (1987);

M. Pospelov et al., Nucl. Rev. D  
78, 115012 (2008);

A. Derevianko et al., Phys. Rev. D  
82, 065006 (2010)

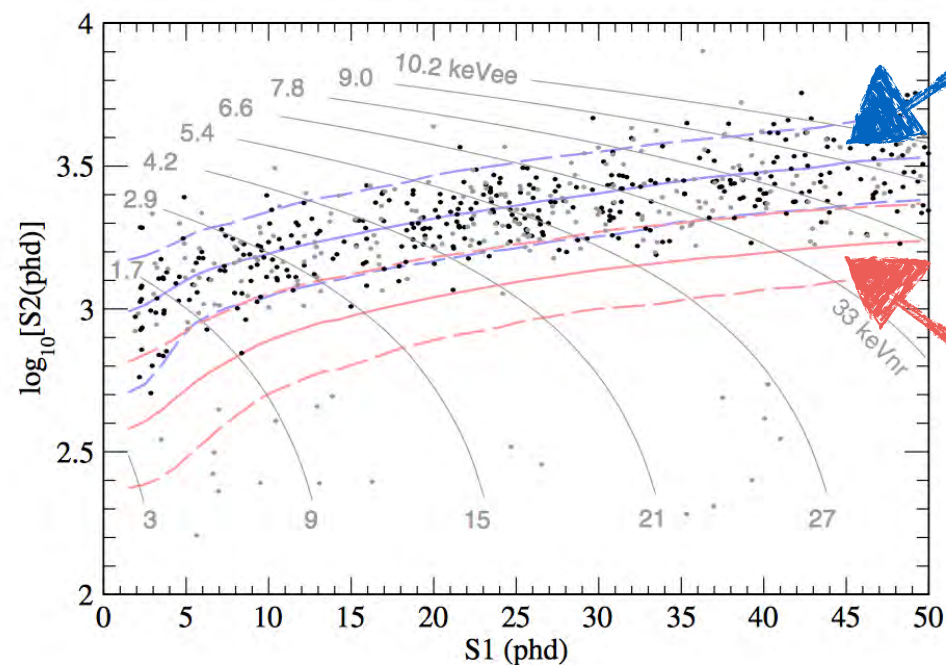


# The use of axio-electric effect to detect axions/ALPs with a xenon TPC



- Detection principle for a TPC: the particle, interacting within the detector, produces photons and electrons
- S1 is the primary scintillation signal (prompt photons)
- S2 is the secondary ionisation signal from electroluminescence (electrons drift thanks to the applied field)

A good discrimination of nuclear (NR) vs electronic recoil (ER) is possible thanks to the ratio  $S2/S1$ : powerful in standard WIMP search, not in axion search



**ER band: most of the background + potential axion signal**

**NR band: few background events + potential WIMP signal**

FIG. 2. Observed events in the 2013 LUX exposure of 95 live days and 145 kg fiducial mass. Points at  $< 18$  cm radius are black; those at  $18$ – $20$  cm are gray. Distributions of uniform-in-energy electron recoils (blue) and an example  $50 \text{ GeV } c^{-2}$  WIMP signal (red) are indicated by 50th (solid), 10th, and 90th (dashed) percentiles of  $S2$  at given  $S1$ . Gray lines, with ER scale of keVee at top and Lindhard-model NR scale of keVnr at bottom, are contours of the linear combined  $S1$ -and- $S2$  energy estimator [19].



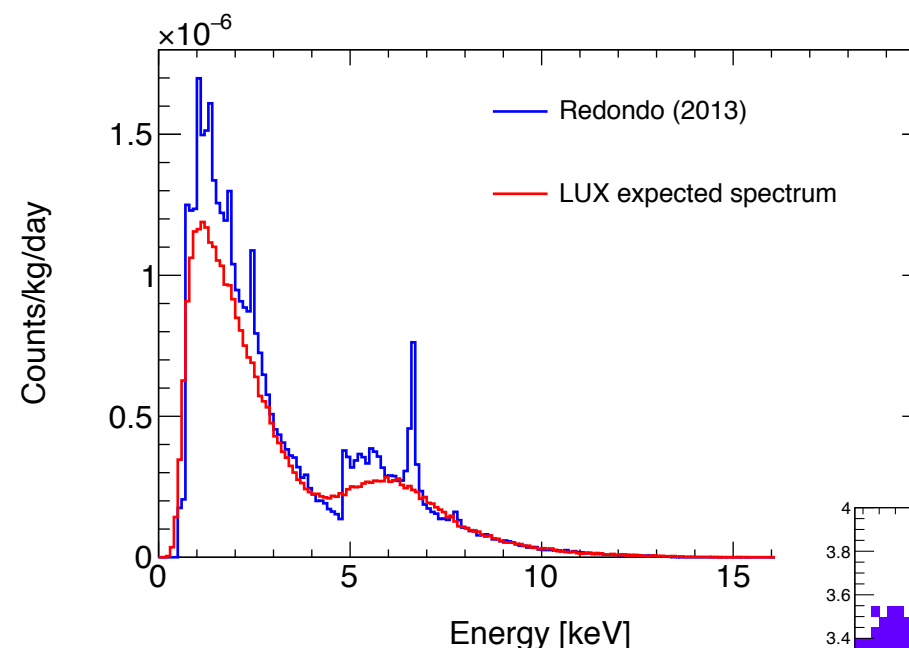
**Axion Results with**



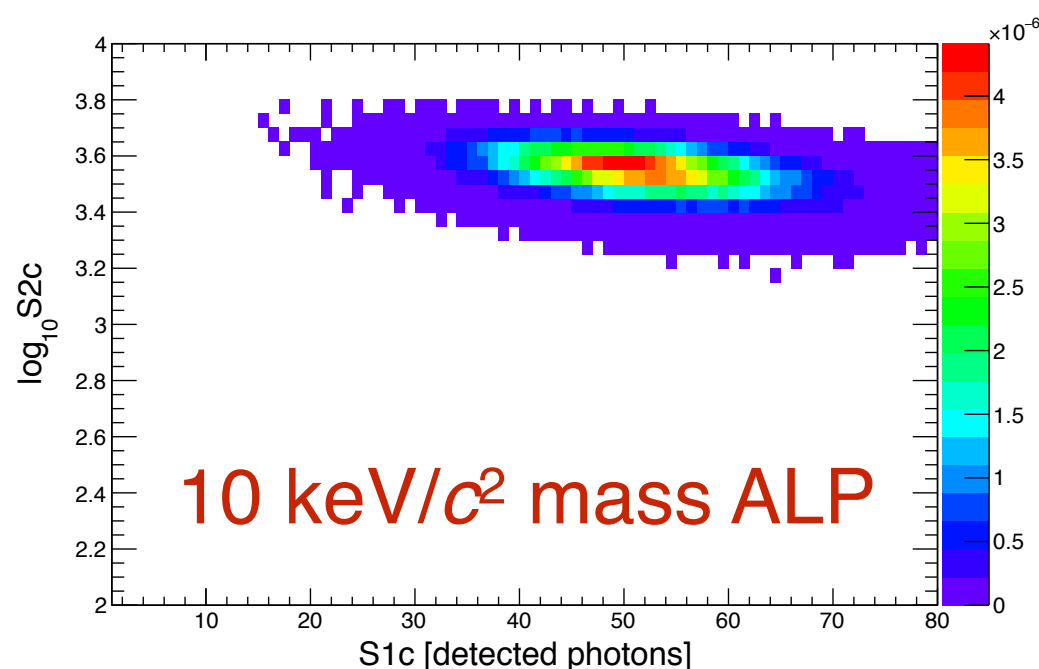
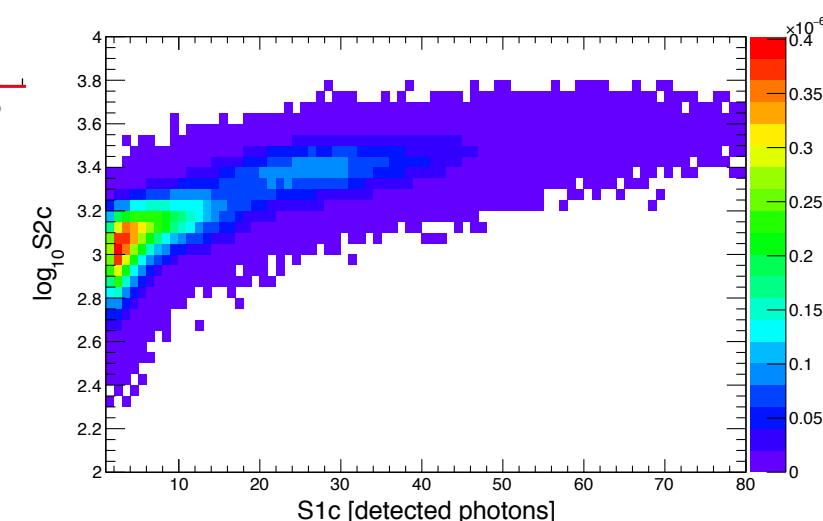


# The signal model

- Solar axion spectral shape: product of solar axion flux [J. Redondo, JCAP 12, 008 (2013)] and photo-electric cross section on xenon, assuming massless axions (still valid for masses smaller than  $1 \text{ keV}/c^2$ )
- Resolution and efficiency effects modelled in accordance with the Noble Element Simulation Technique (NEST) package [M. Szydagis et al., JINST 6, P10002 (2011)]



Solar axions



- ALPs expected to be essentially at rest within the galaxy
- Axio-electric absorption leads to electron recoils with kinetic energy equal to the ALP mass: sharp spectral feature, smeared by energy resolution

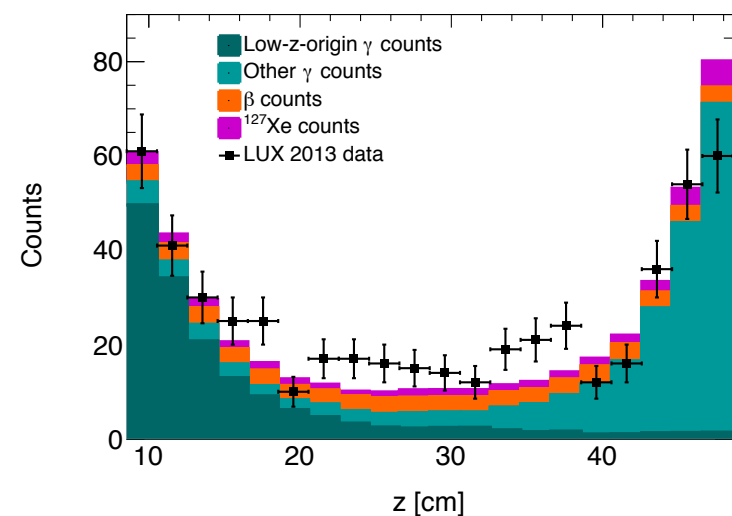
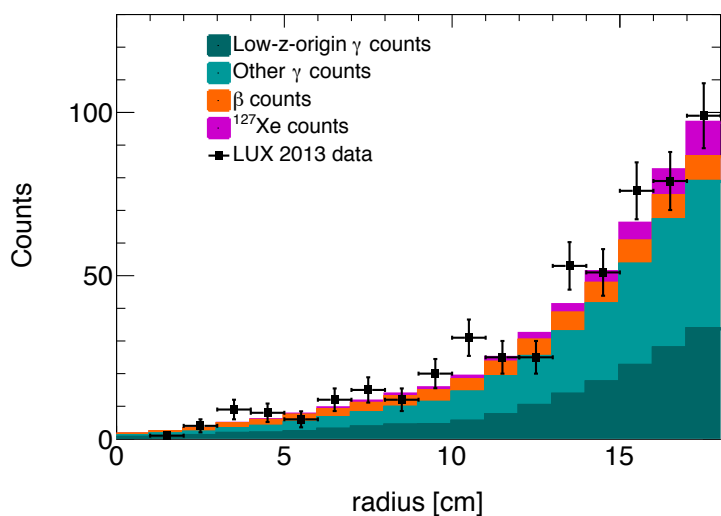
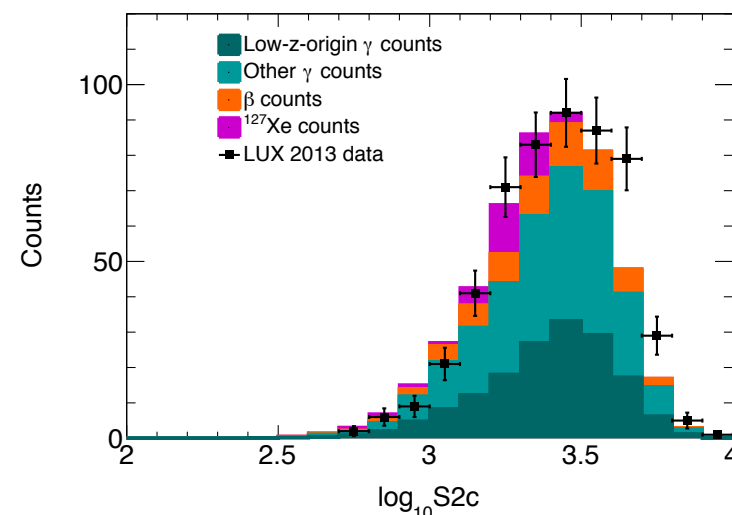
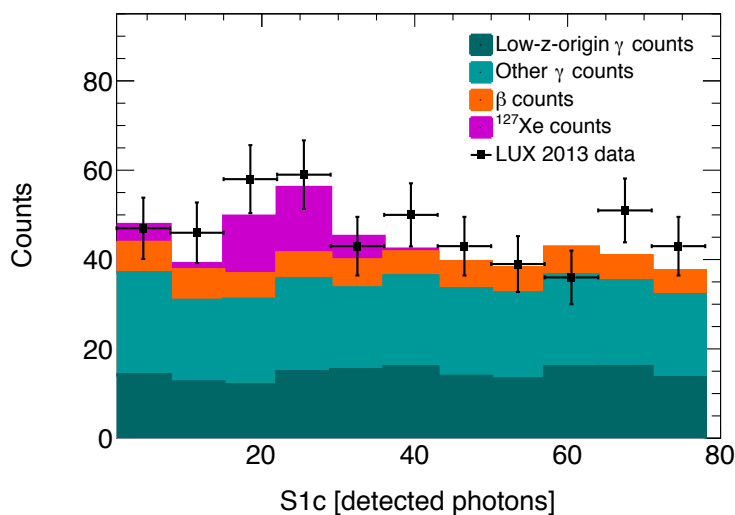
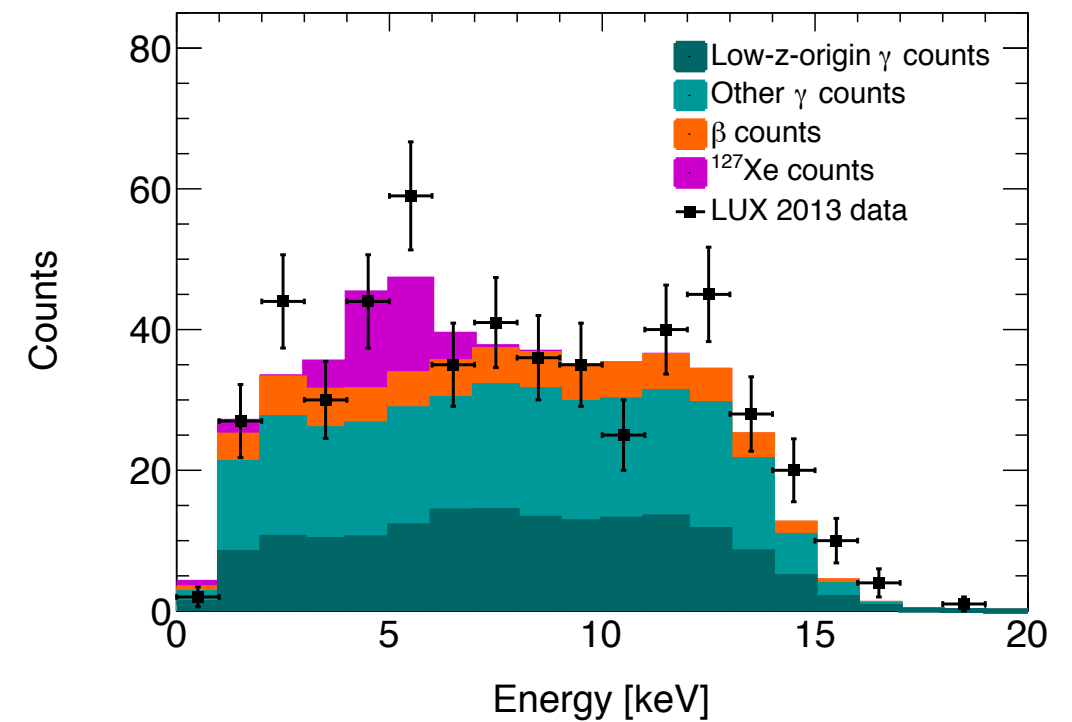


# LUX 2013 search data & the background model

- LUX 2013 data: 95 live-days, 118 kg fiducial mass
- Low rate of background radioactivity thanks to detector design, location deep underground, construction materials, xenon self-shielding, active circulation and purification
- Different contributions to the background:
  - Compton scattering of gamma rays from detector component radioactivity
  - $^{85}\text{Kr}$  and Rn-daughter contaminants undergoing beta decays
  - x-rays emitted following those  $^{127}\text{Xe}$  electron-capture decays where the coincident gamma escapes the xenon



- LUX 2013 data and background model as a function of recoil energy, with the energy reconstructed as  $E = (S1c/g1 + S2c/(\epsilon g2))W$
- $g1$ : geometric light collection efficiency and PMT quantum efficiency
- $\epsilon g2$ : electron extraction efficiency and number of photons detected per electron extracted



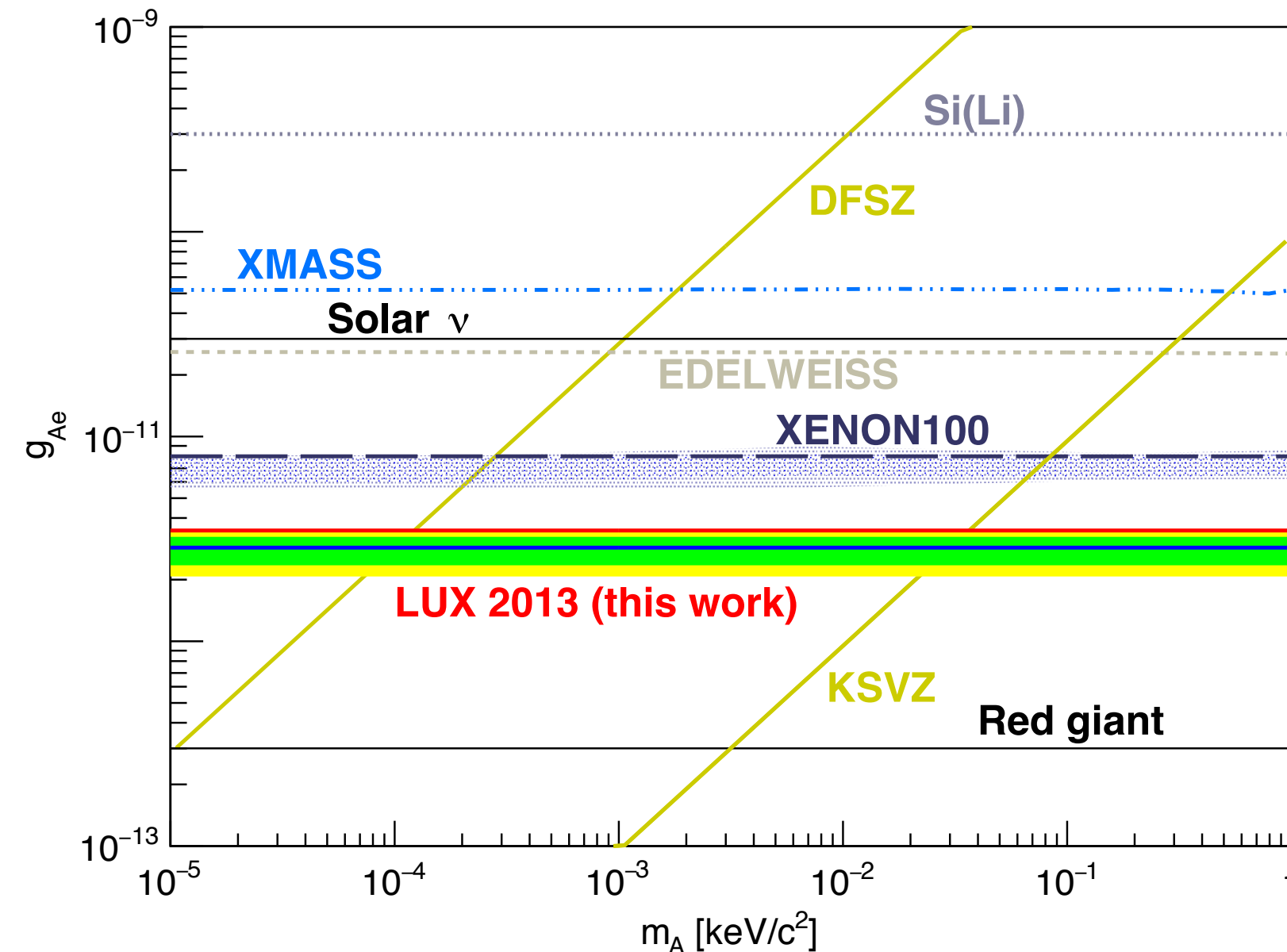
- Backgrounds modelled on the four observables used in the statistical analysis: the prompt scintillation (S1) and the logarithm in base 10 of the proportional (S2) signal, and the radius (r) and depth (z) of the event location

- Statistical PLR (Profile Likelihood Ratio) analysis, aimed at setting a two-sided limit on the coupling between axions/ALPs and electrons  $g_{Ae}$ , having the BG rates as nuisance parameters



# solar axions result

arXiv:1704.02297 [astro-ph.CO] (2017)



- QCD axion theoretical models:
  - DFSZ**: axion is the phase of a new electroweak singlet scalar field and couples to a new heavy quark, not to SM ones
  - KSVZ**: axion does not couple directly to quarks and leptons, but via its interaction with two Higgs doublets
- Red Giant** limit: the degenerate core of a low-mass red giant before helium ignition is a helium white dwarf; the observed white-dwarf luminosity function reveals that their cooling speed agrees with expectations, constraining new cooling agents such as axion emission

LUX 2013 excludes  $g_{Ae} > 3.5 \times 10^{-12}$  (90% CL)

LUX 2013 excludes  $m_A > 0.12 \text{ eV}/c^2$  (DFSZ model)

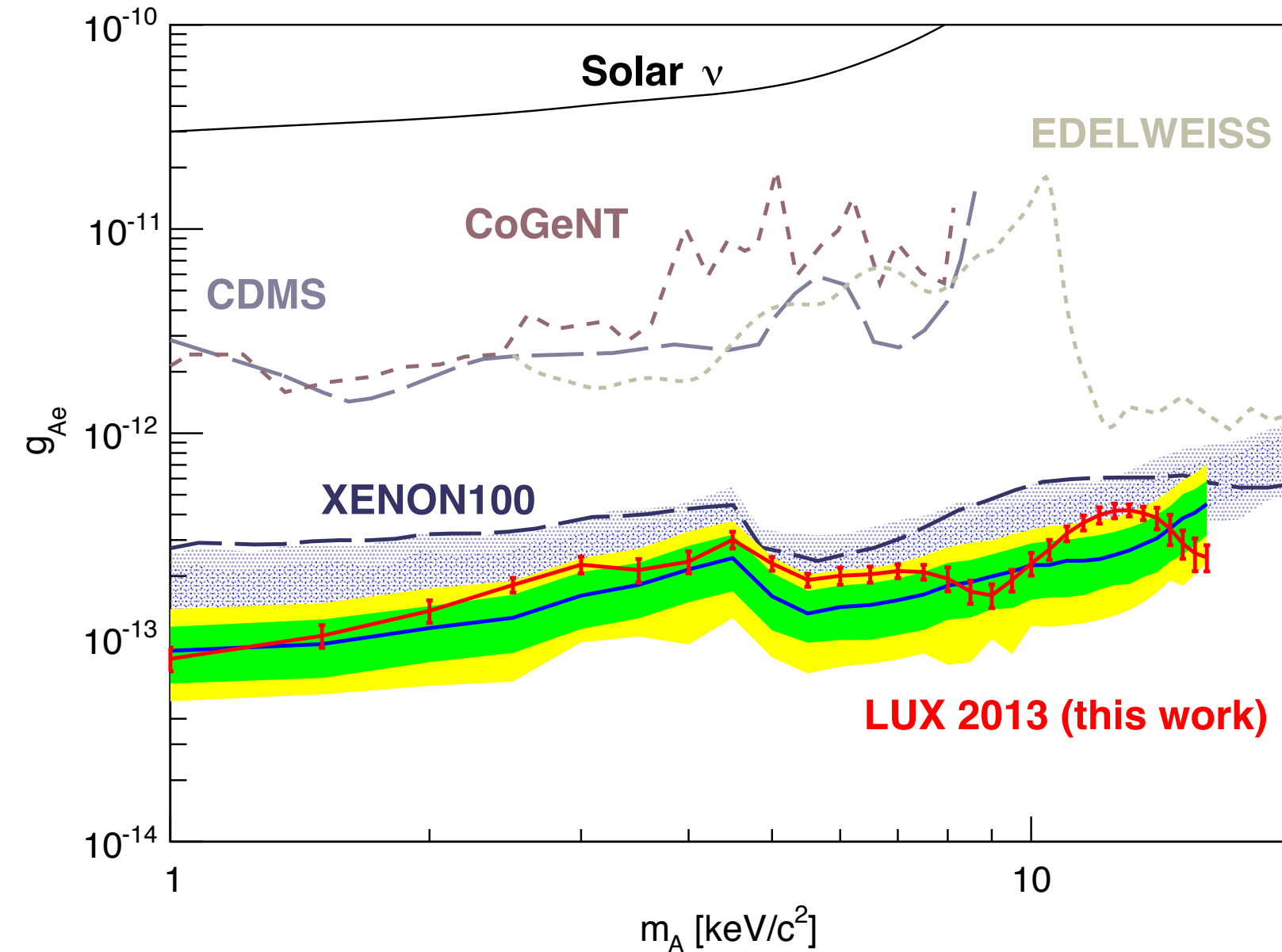
LUX 2013 excludes  $m_A > 36.6 \text{ eV}/c^2$  (KSVZ model)





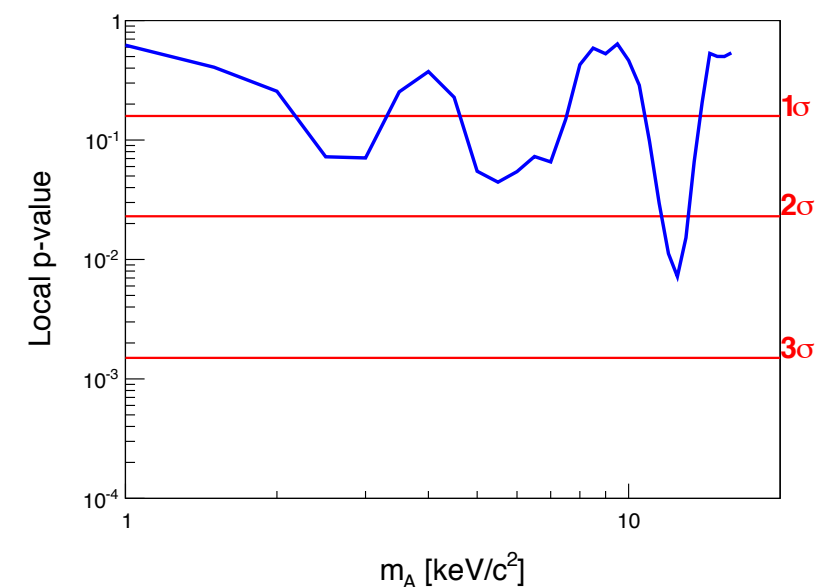
# galactic ALPs result

arXiv:1704.02297 [astro-ph.CO] (2017)



LUX 2013 excludes  $g_{Ae} > 4.2 \times 10^{-13}$  (90% CL) across the range 1-16 keV/c<sup>2</sup> in ALP mass

- Local p-value as a function of the ALP mass, with the corresponding number of standard deviations ( $\sigma$ ) away from the null hypothesis
- At 12.5 keV/c<sup>2</sup> a local p-value of  $7.2 \times 10^{-3}$  (2.4 $\sigma$  deviation) corresponds to a global p-value of  $5.2 \times 10^{-2}$  (1.6 $\sigma$  deviation) — applying the Look Elsewhere Effect [E. Gross and O. Vitells, Eur. Phys. J., C70:525 (2010)]



# Future Searches with





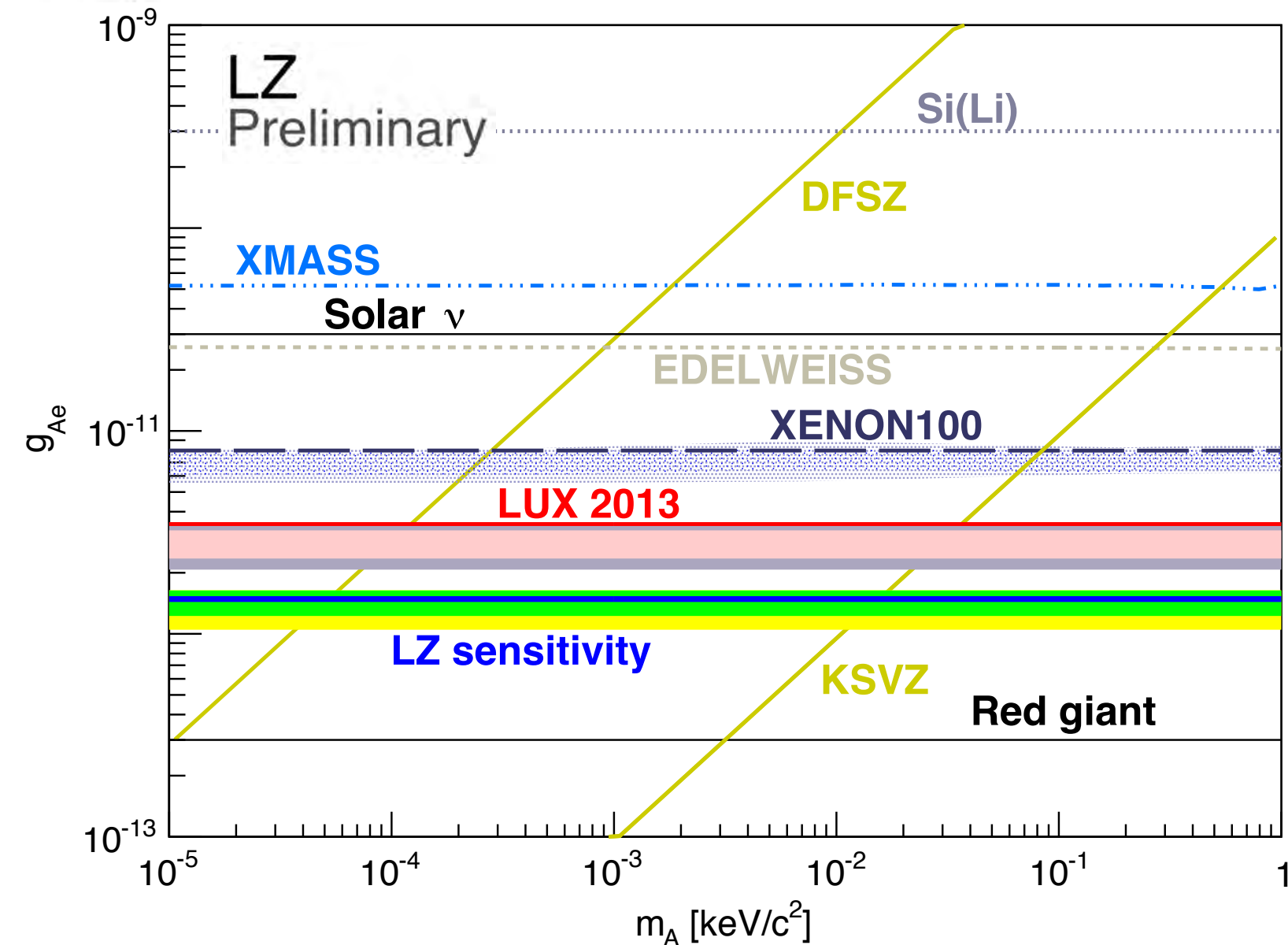


# Sensitivity for axion searches with LZ

- LZ expected exposure: 1000 live-days, 5.6 ton fiducial mass
- Low rate of background radioactivity thanks to detector design, location deep underground, construction materials and xenon self-shielding
- Different contributions to the background:
  - Compton scattering of gamma rays from detector component radioactivity
  - $^{85}\text{Kr}$  and Rn-daughter contaminants undergoing beta decays
  - neutrinos (atmospheric, PP solar)
- Statistical PLR (Profile Likelihood Ratio) analysis, aimed at setting a two-sided limit on the coupling between axions/ALPs and electrons  $g_{\text{Ae}}$ , having the BG rates as nuisance parameters and S1 and S2 as observables



# most recent solar axions sensitivity projection



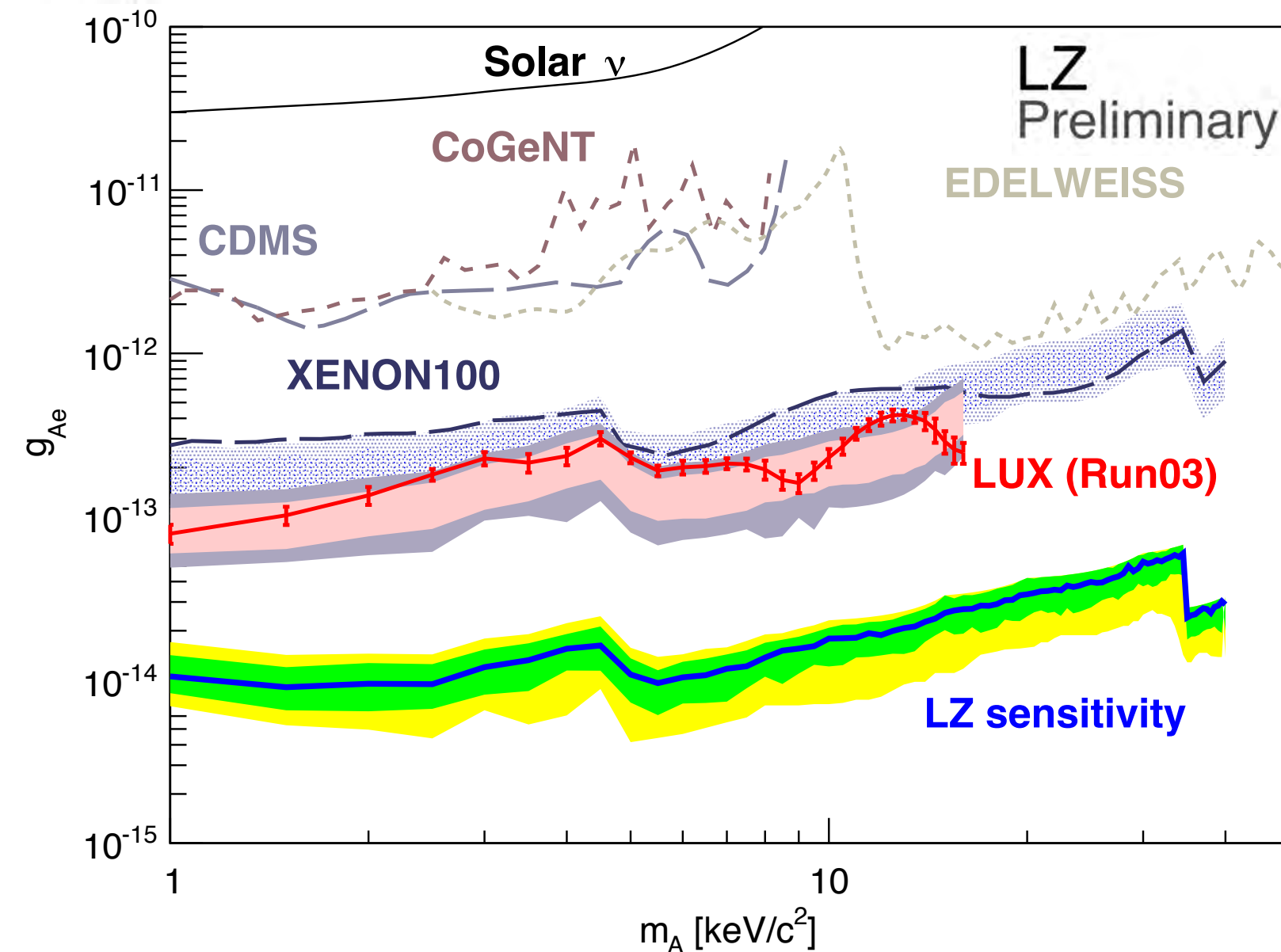
LZ sensitivity (1000 live-days, 5.6 ton fiducial mass)  
excludes  $g_{Ae} > 1.5 \times 10^{-12}$  (90% CL)

- QCD axion theoretical models:
  - DFSZ**: axion is the phase of a new electroweak singlet scalar field and couples to a new heavy quark, not to SM ones
  - KSVZ**: axion does not couple directly to quarks and leptons, but via its interaction with two Higgs doublets
- Red Giant** limit: the degenerate core of a low-mass red giant before helium ignition is a helium white dwarf; the observed white-dwarf luminosity function reveals that their cooling speed agrees with expectations, constraining new cooling agents such as axion emission





# most recent galactic ALPs sensitivity projection



LZ sensitivity (1000 live-days, 5.6 ton fiducial mass)  
excludes  $g_{Ae} > 5.9 \times 10^{-14}$  (90% CL)  
across the range 1-40 keV/c<sup>2</sup> in ALP mass

# Summary

- Xenon detectors (such as LUX and LZ) present suitable characteristics to test models beyond the standard WIMP scenario
- QCD axions can solve the strong CPV problem
- Some classes of axions/axion-like particles are also suitable Dark Matter candidates
- Axion signal in a xenon TPC is expected in the ER band, where also most of the backgrounds sit
- A Profile Likelihood Ratio statistical strategy is used to set an upper limit on the coupling  $g_{Ae}$ , using the most meaningful experimental quantities as observables
- LUX 2013 delivers the most stringent constraints to date for these interactions (**yesterday on arXiv:1704.02297**) and the analysis of the full LUX exposure is planned
- LZ will be able to probe a wider phase space thanks to its sensitivity



# Thank you!

It's been a very interesting year for LUX!

arXiv:1704.02297 (2017)

Phys. Rev. D 95, 012008 (2017)

Phys. Rev. Lett. 118, 021303 (2017)

Phys. Rev. Lett. 116, 161302 (2016)

Phys. Rev. Lett. 116, 161301 (2016)

Phys. Rev. D 93, 072009 (2016)

## The LUX collaboration

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Dana Byram	Support Scientist

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Paul Terman	Graduate Student

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Scott Hertel	Postdoc
Kevin O'Sullivan	Postdoc
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Evan Pease	Graduate Student
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Nicole Larsen	Graduate Student

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Aaron Manalaysay	Project Scientist
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James Morad	Graduate Student
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Dean White	Engineer
Carmen Carmona	Postdoc
Scott Haselschwardt	Graduate Student
Curt Nehr Korn	Graduate Student
Melih Solmaz	Graduate Student

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Shaun Alsum	Graduate Student

## The Z collaboration

31 institutions  
~200 people



LIP Coimbra (Portugal)

Center for Underground Physics (Korea)

MEPhI (Russia)

Edinburgh University (UK)

University of Liverpool (UK)

Imperial College London (UK)

University College London (UK)

University of Oxford (UK)

STFC Rutherford Appleton Laboratories (UK)

University of Sheffield (UK)

University of Alabama

University at Albany SUNY

Berkeley Lab (LBNL)

University of California, Berkeley

Brookhaven National Laboratory

Brown University

University of California, Davis

Fermi National Accelerator Laboratory

Lawrence Livermore National Laboratory

University of Maryland

University of Michigan

Northwestern University

University of Rochester

University of California, Santa Barbara

University of South Dakota

South Dakota School of Mines & Technology

South Dakota Science and Technology Authority

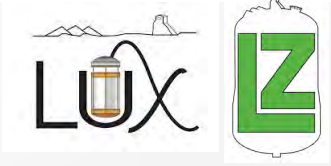
SLAC National Accelerator Laboratory

Texas A&M

Washington University

University of Wisconsin





# *Back-up slides*



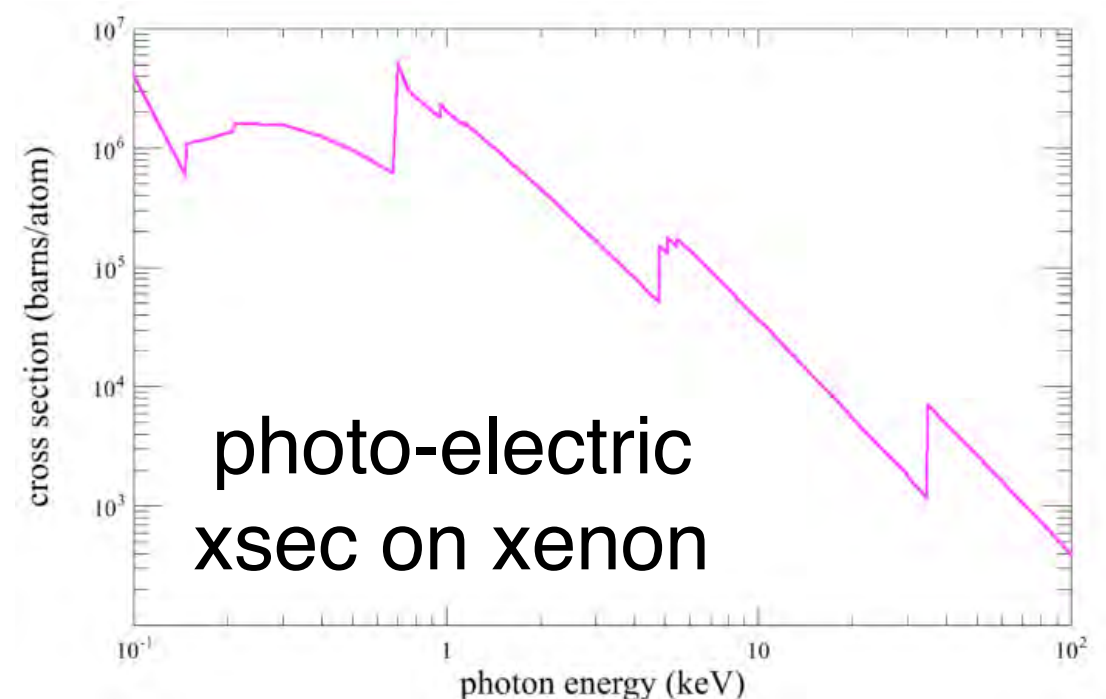
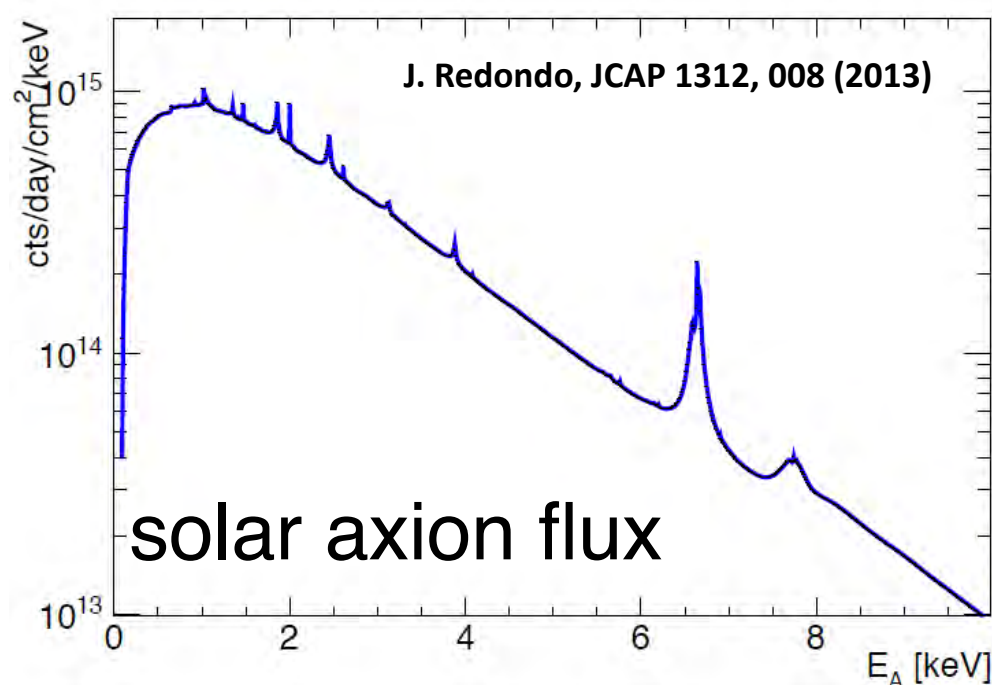
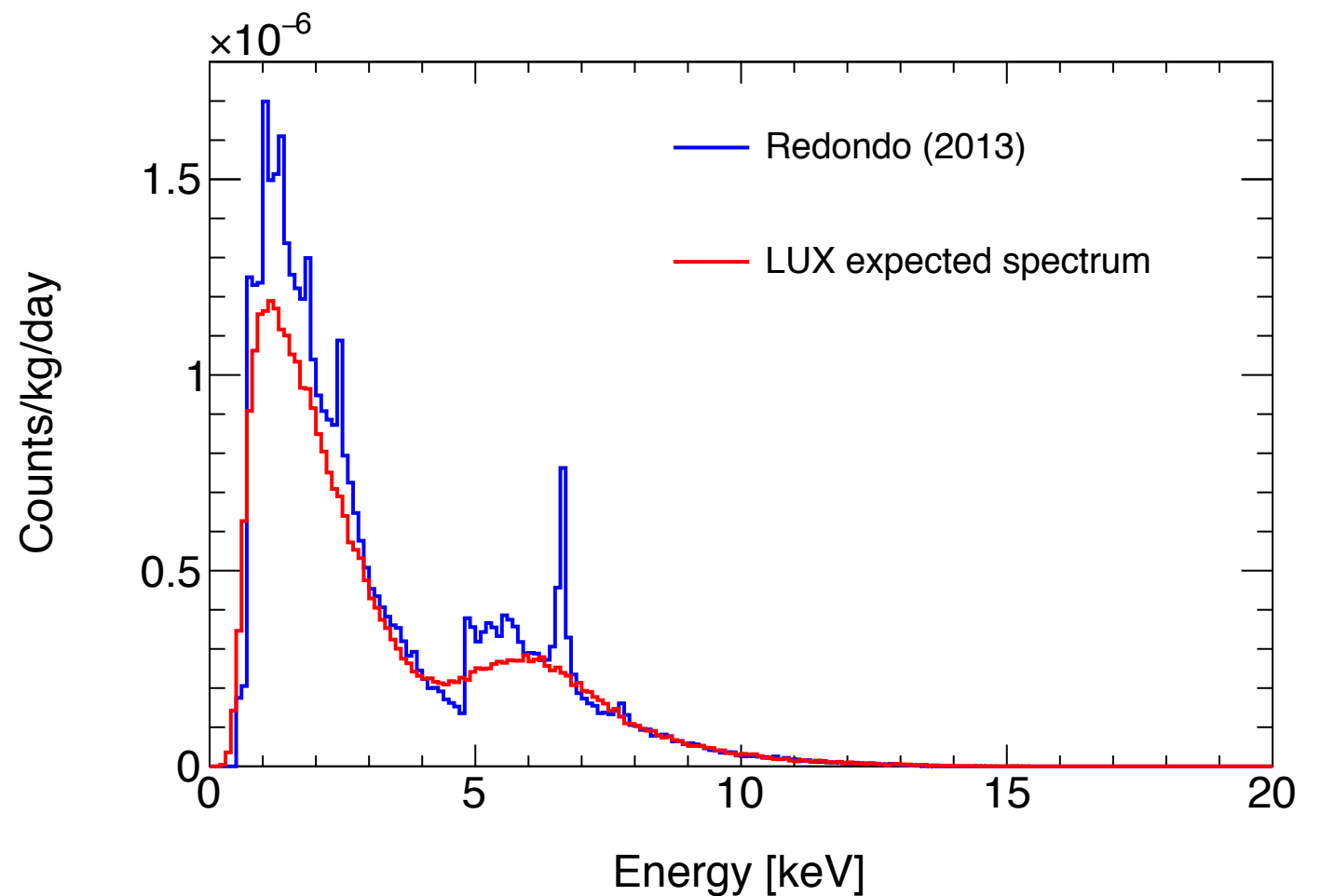


# Limit conversion: nSig to gAe

- Limit on  **$gAe = gAe_{sim} * (nSig/nPDF)^{power}$** 
  - $gAe_{sim}$  = arbitrary coupling used to generate the signal model
  - nSig = limit set by the PLR on the number of events
  - nPDF = integral of the signal PDFs \* exposure \* fiducial mass
  - power varies with the axion type
    - it is 0.25 for solar axions, as the interaction rate scales with  $gAe^4$
    - it is 0.50 for galactic ALPs, as the interaction rate scales with  $gAe^2$

# The solar axions signal model

- Solar axion spectral shape: product of solar axion flux [J. Redondo, JCAP 12, 008 (2013)] and photo-electric cross section on xenon, assuming massless axions (still valid for masses smaller than  $1 \text{ keV}/c^2$ )
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# LUX efficiency for electronic recoils

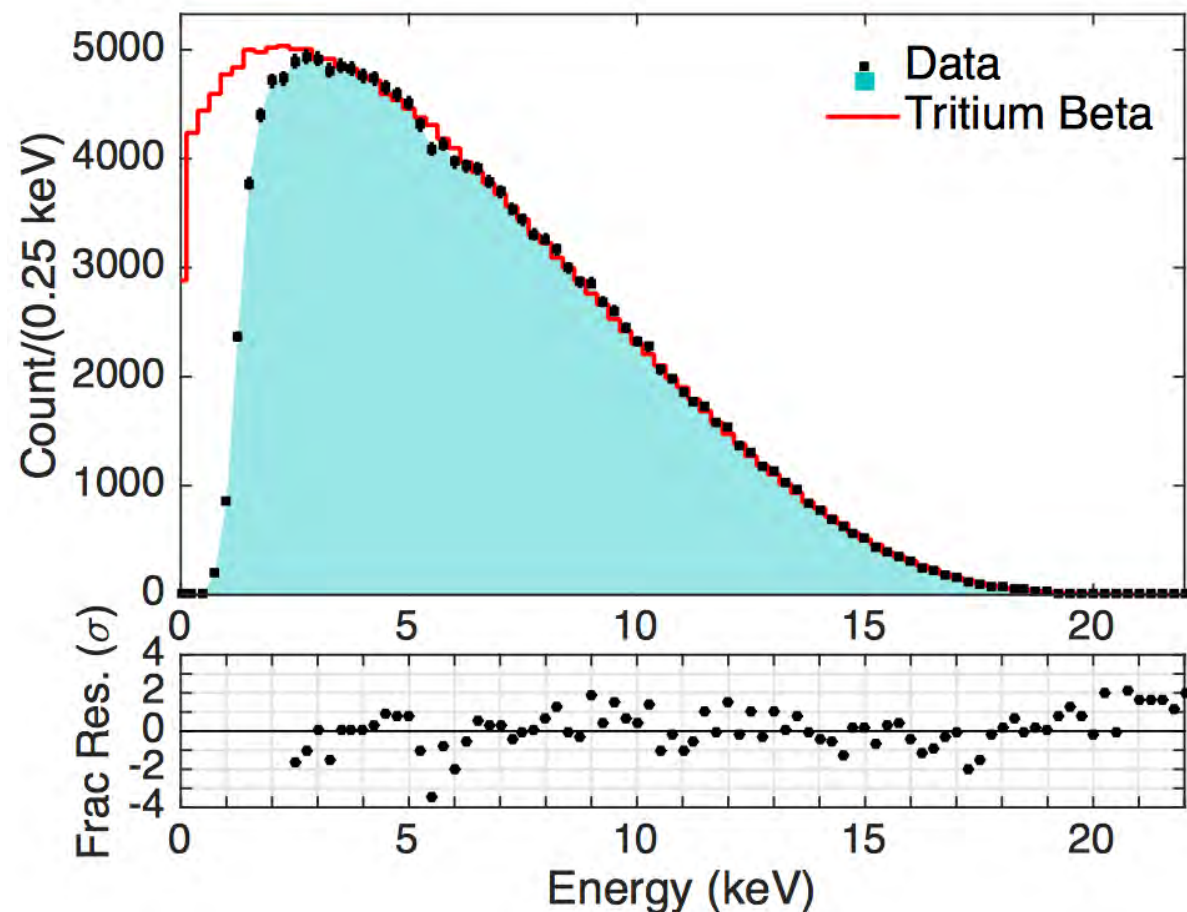


FIG. 5: Top: The tritium energy spectrum measured by LUX with the combined energy model (black) compared to a tritium spectrum convolved with detector resolution ( $\frac{\sigma_E}{W} = \sqrt{\sigma^2(n_\gamma) + \sigma^2(n_e)}$ ). The p-value between data and model from 3 to 18 keV is 0.70. Bottom: Bin-by-bin fit residuals between data and theory, in units of  $\sigma$ .

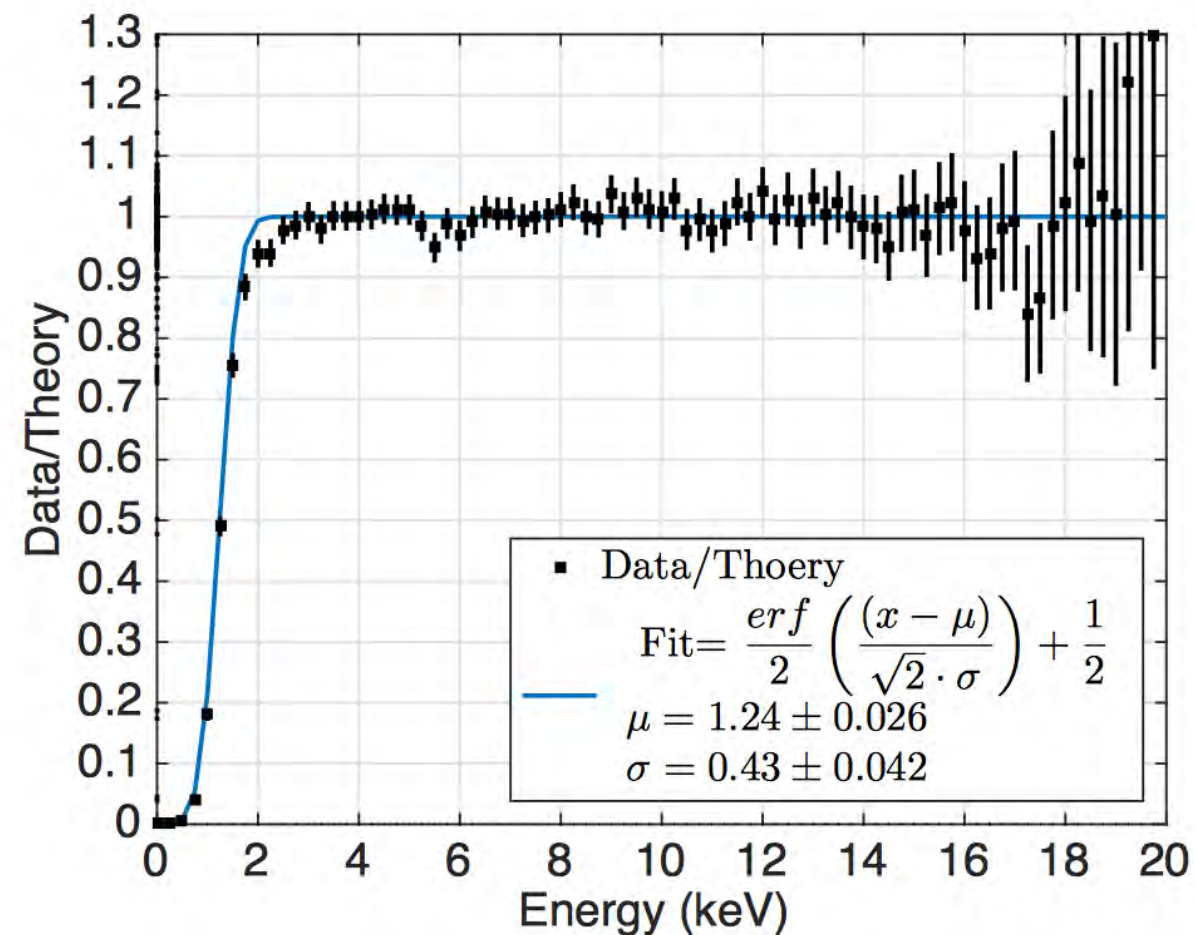


FIG. 6: Ratio of the measured tritium energy spectrum and the true one convolved with the detector resolution. A fit to an error function is shown.

D. S. Akerib et al., Phys. Rev. D93, 072009 (2016)



# LUX energy resolution for electronic recoils

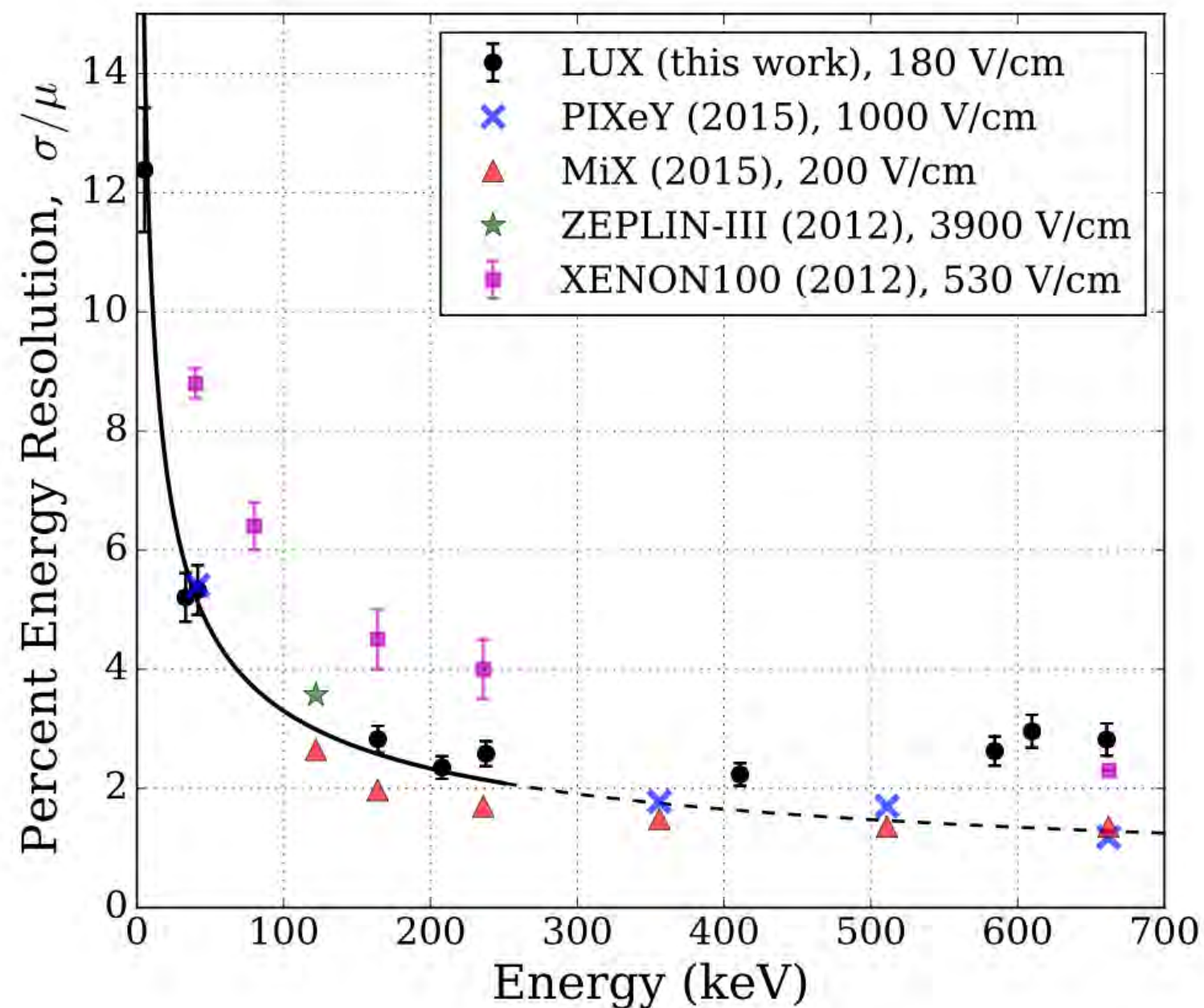
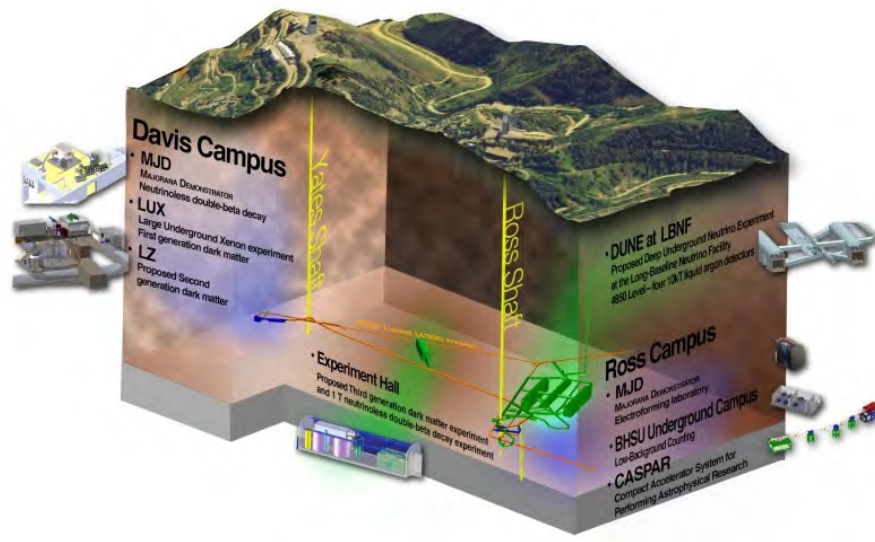


FIG. 8. The measured energy resolution at known energy peaks in the LUX ER backgrounds. The detector is optimized for low energy sensitivity, and variable amounts of PMT saturation and single-electron contributions affect S2 pulses and hamper the energy resolution at high energy, as discussed in the text. Data from the PIXeY (blue x; [26, 27]), MiX (red triangle; [28]), ZEPLIN-III (green star; [29]), and XENON100 (magenta square; [30]) are shown for comparison.

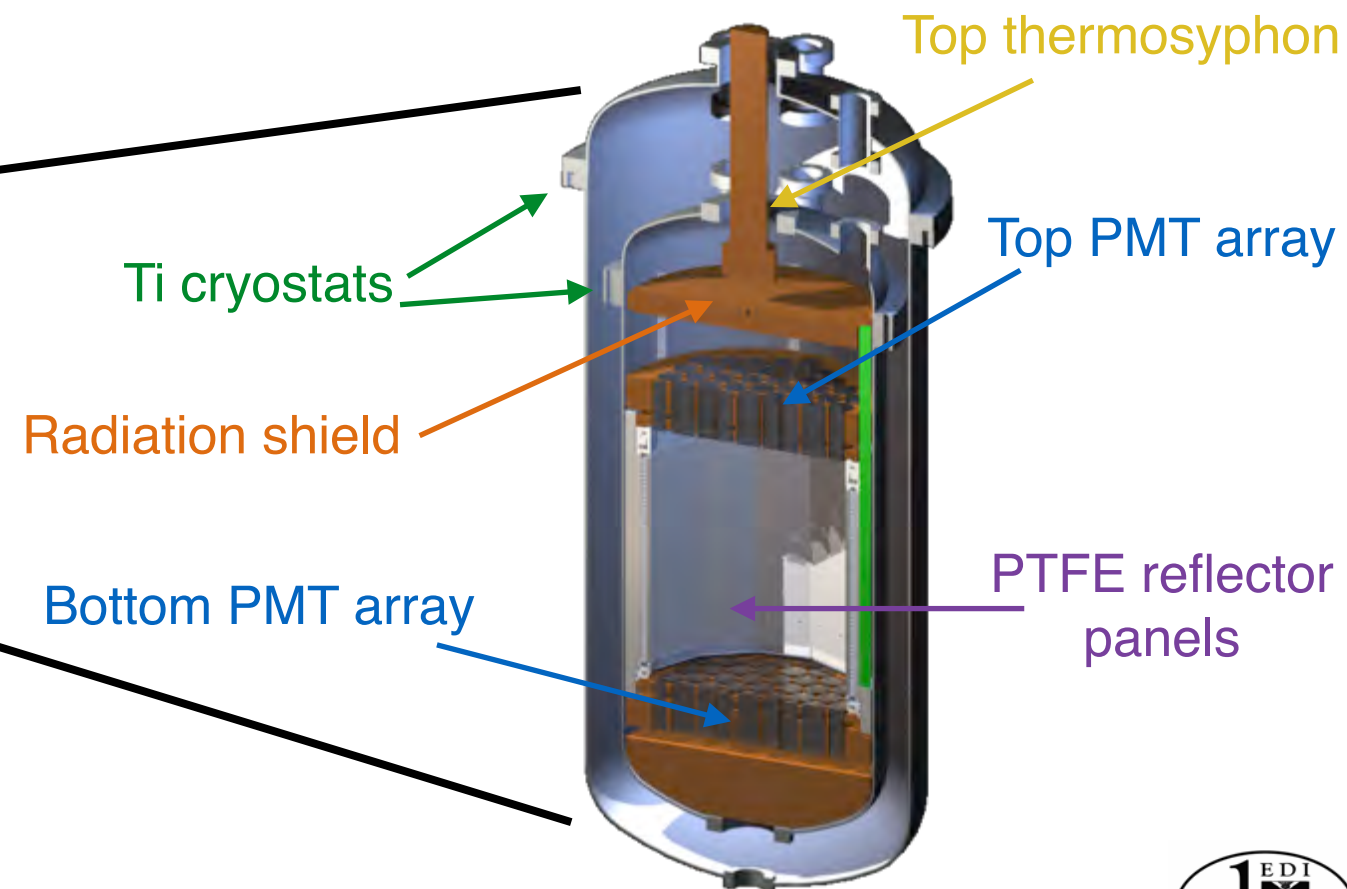
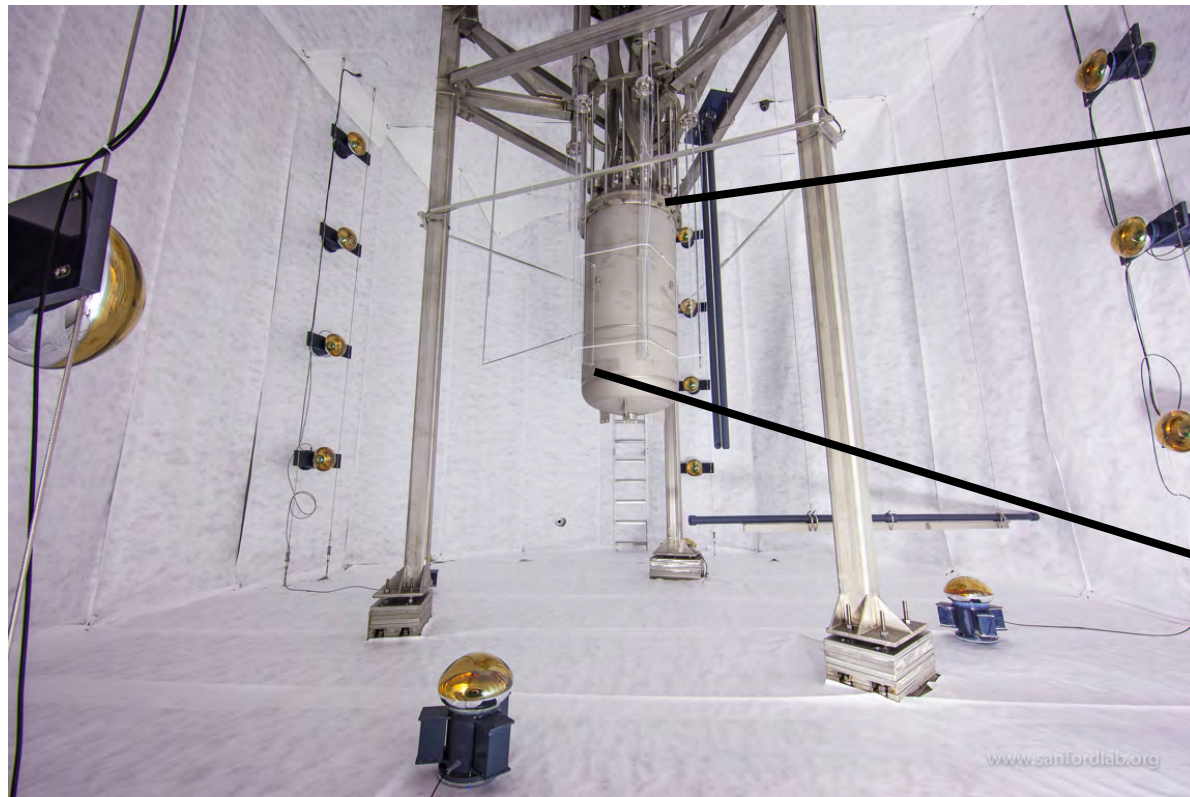
D. S. Akerib et al., Phys. Rev. D95, 012008 (2017)

# The Large Underground Xenon experiment



- 370 kg of liquid xenon, 250 kg of active mass
- with a layer of gaseous xenon maintained above the liquid xenon (dual phase TPC)
- Vertical electric field applied (181 V/cm)
- 61 top + 61 bottom PMTs to detect signals

- LUX used to operate 4850 feet underground, in Davis Carven of the Sanford Underground Research Facility (South Dakota, USA)

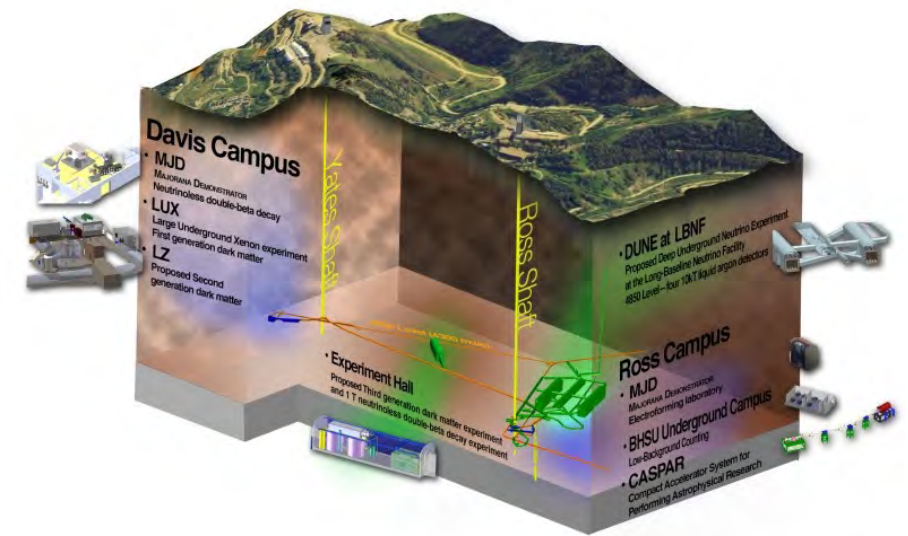
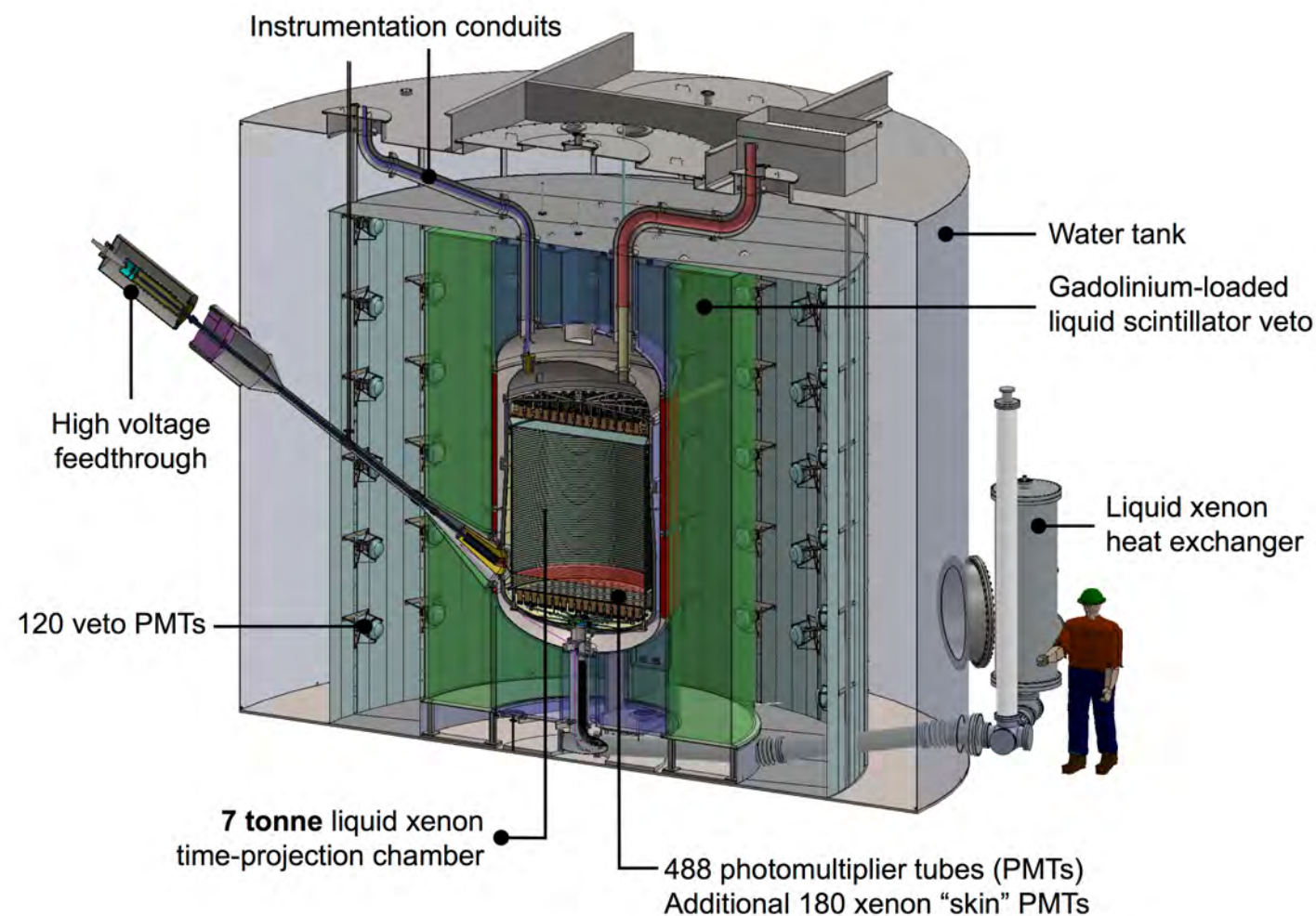




# The LZ (LUX+ZEPLIN) experiment

- Dual-phase xenon TPC: 10 ton total mass, 7 ton active LXe mass, 5.6 ton fiducial mass

The LZ Dark Matter Experiment



- Will be installed at SURF, where LUX used to work — onsite improvements in infrastructure for LZ

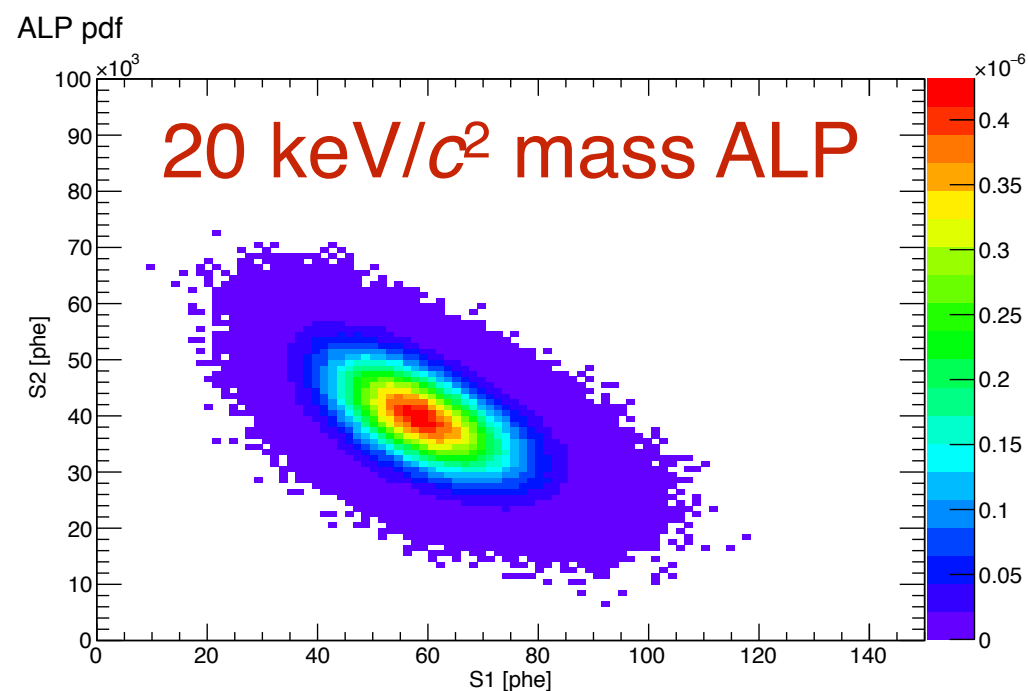
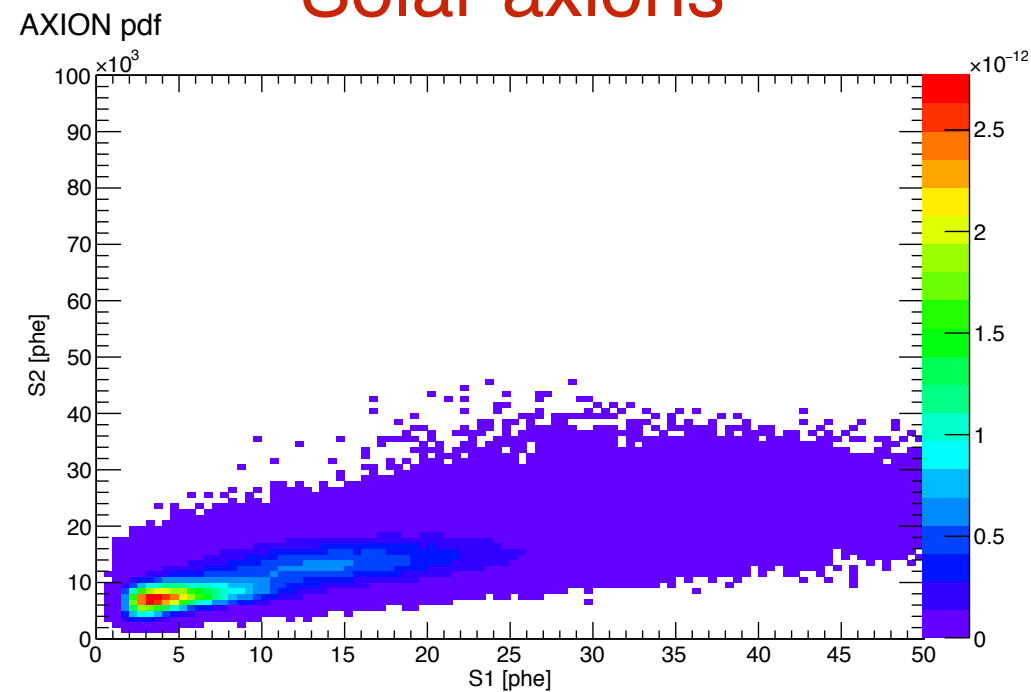




# The signal model

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## Solar axions



- ALPs expected to be essentially at rest within the galaxy
- Axio-electric absorption leads to electron recoils with kinetic energy equal to the ALP mass: sharp spectral feature, smeared by energy resolution

# Why axions ?

## (Particle Physics)

- The Strong CP violation problem

- the QCD Lagrangian acquires a term, proportional to a static parameter  $\theta$ , because of the non zero divergence of the axial current

- this term is CP violating, but we do not observe any CP violation in strong interactions

$$L_{QCD} = \bar{\psi}(i\gamma^\mu D_\mu - m)\psi - \frac{1}{4}G^{a\mu\nu}G_{\mu\nu}^a + \frac{\alpha_S\bar{\theta}}{8\pi}G_{\mu\nu}^a\hat{G}^{a\mu\nu}$$

- The Peccei and Quinn solution (1977)

- a new global symmetry  $U(1)_{PQ}$  is introduced and spontaneously broken at some large energy scale,  
and the axion is the

Nambu-Goldstone boson generated

$$L = \bar{\psi}(i\gamma^\mu D_\mu - m)\psi - \frac{1}{4}G^{a\mu\nu}G_{\mu\nu}^a - \frac{1}{2}\partial_\mu a_{phys}\partial^\mu a_{phys} + L_{int}[\partial^\mu a_{phys}/f, \psi] + \frac{a_{phys}}{f_a}\xi\frac{\alpha_S}{8\pi}G_{\mu\nu}^a\hat{G}^{a\mu\nu}$$

- the axion field terms introduced in the QCD Lagrangian, cancel out the term proportional to  $\theta$ , providing a dynamical solution to the strong CP problem