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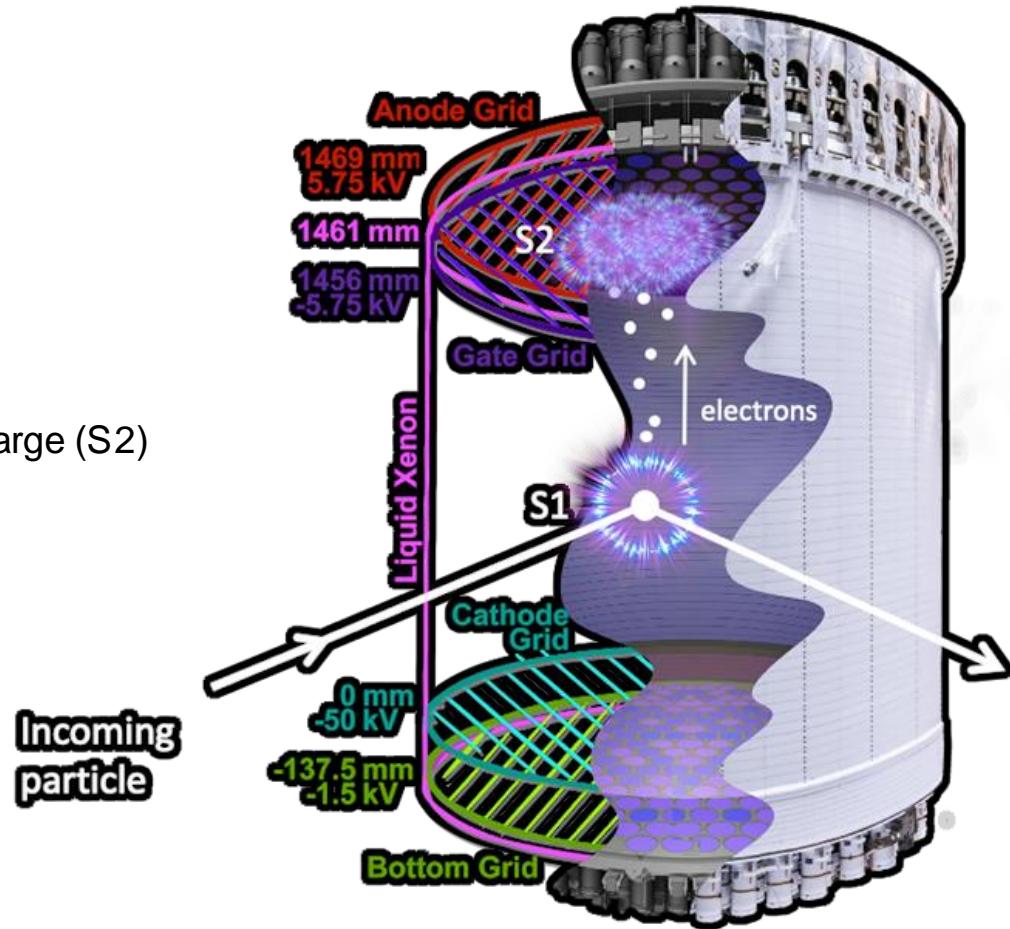
Going beyond SI
searches with
LUX-ZEPLIN

Sam Eriksen, on behalf of the LZ Collaboration
EPS-HEP 2023
25th August 2023

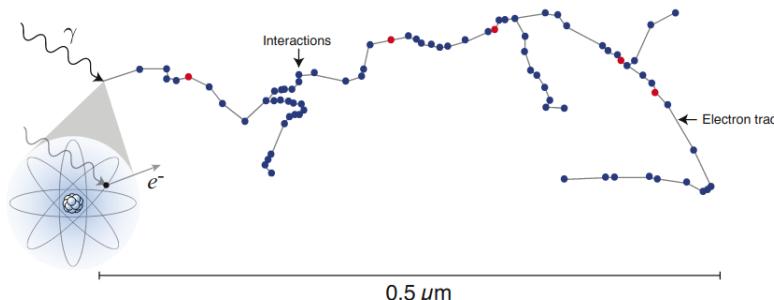


Dual Phase Time Projection Chamber

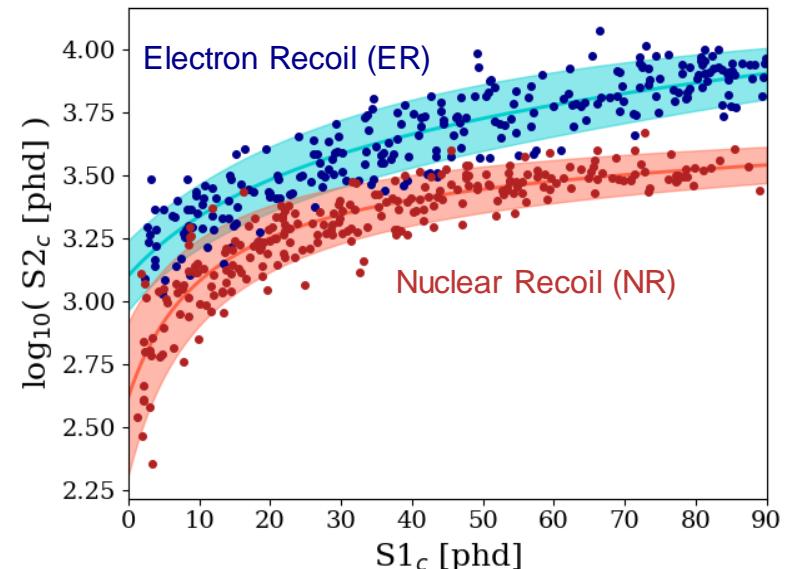
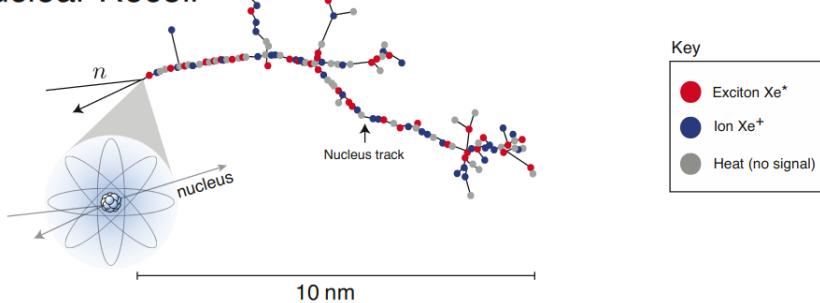
- Primary scintillation light (S1)
- Secondary scintillation induced from free charge (S2)
- 3D reconstruction allows for fiducialisation
- ER/NR discrimination from S1:S2 ratio



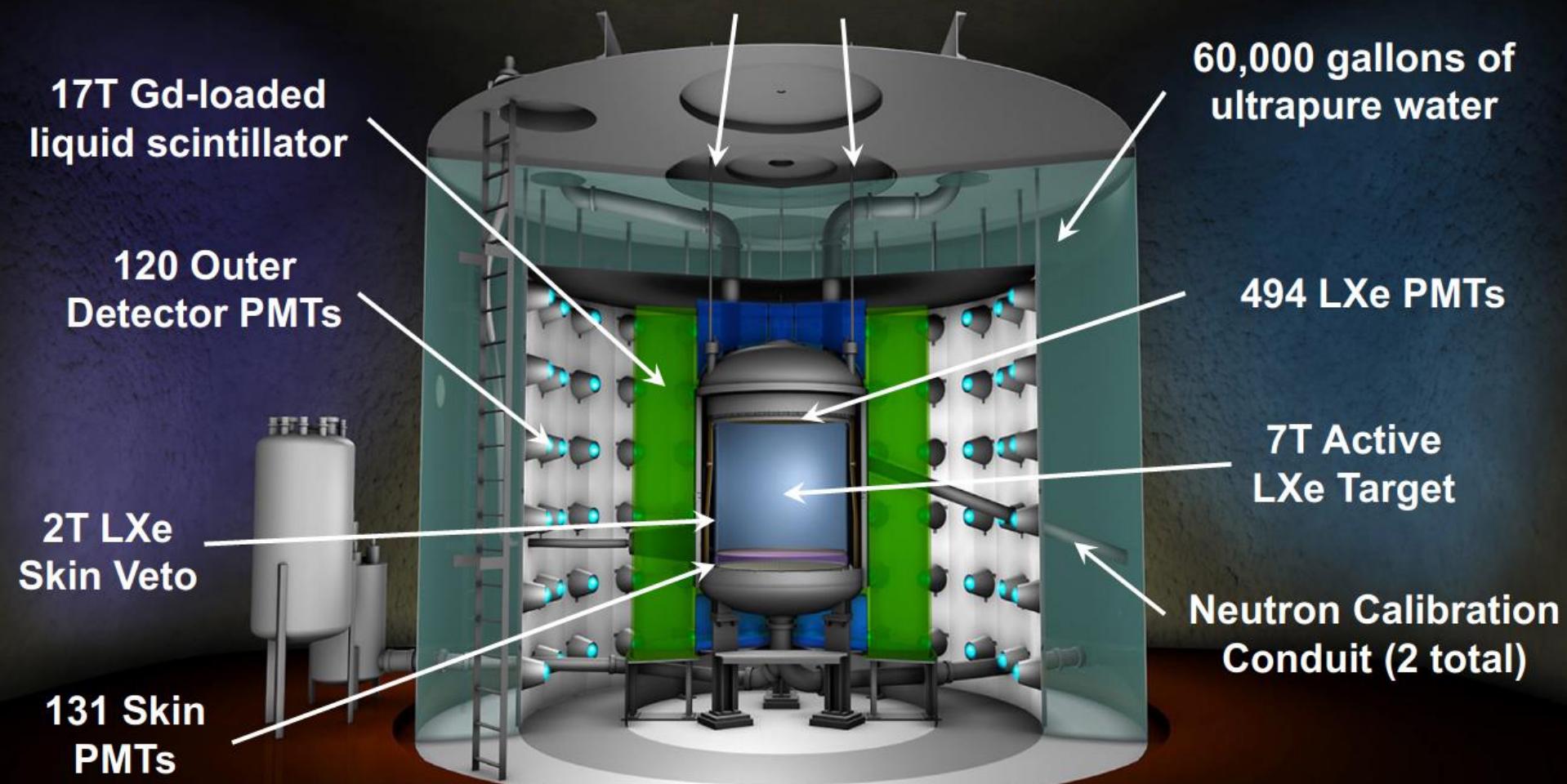
Electron Recoil



Nuclear Recoil



Calibration Source Deployment Tubes (3 Total)





Detector Conditions

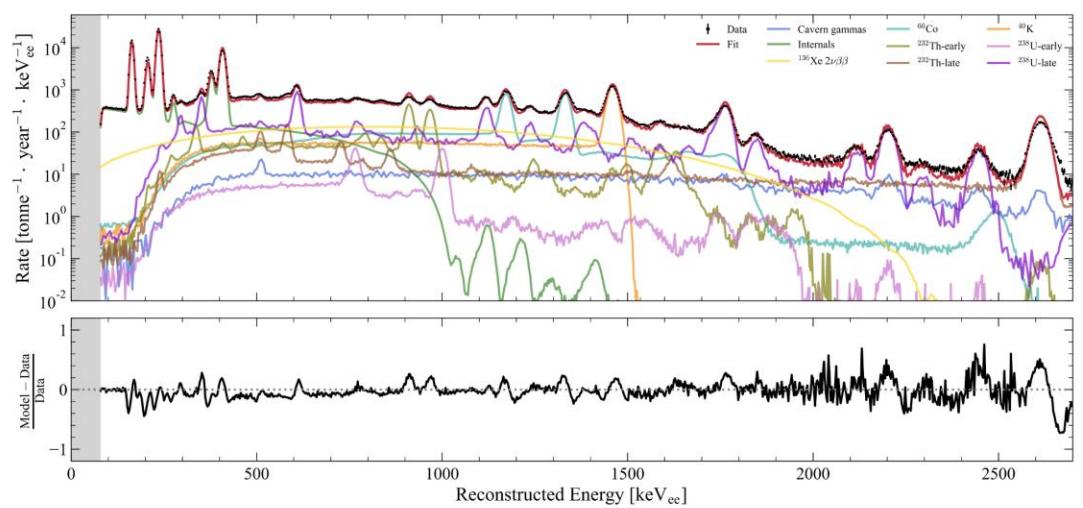
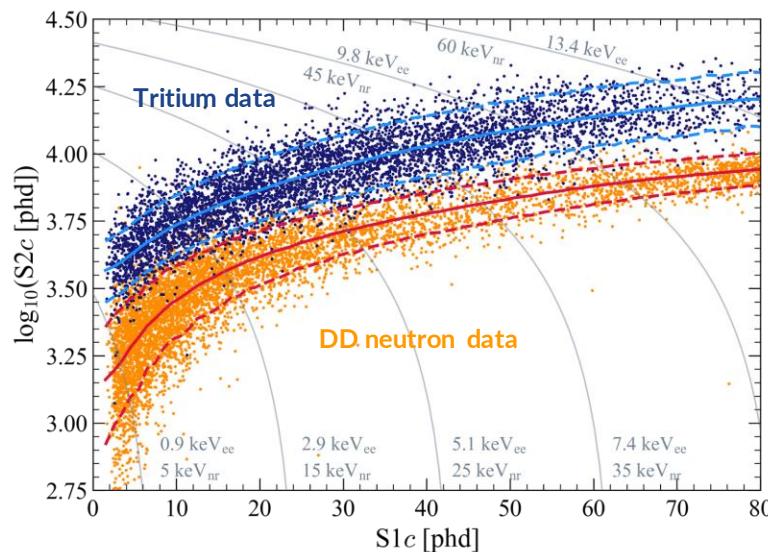
- Drift field: 193 V/cm
- Extraction field: 7.3 kV/cm in gas
- >97% if PMTs operational
- Liquid temperature (174.1 K)
- 3.3 t/day Xe purified through hot getter

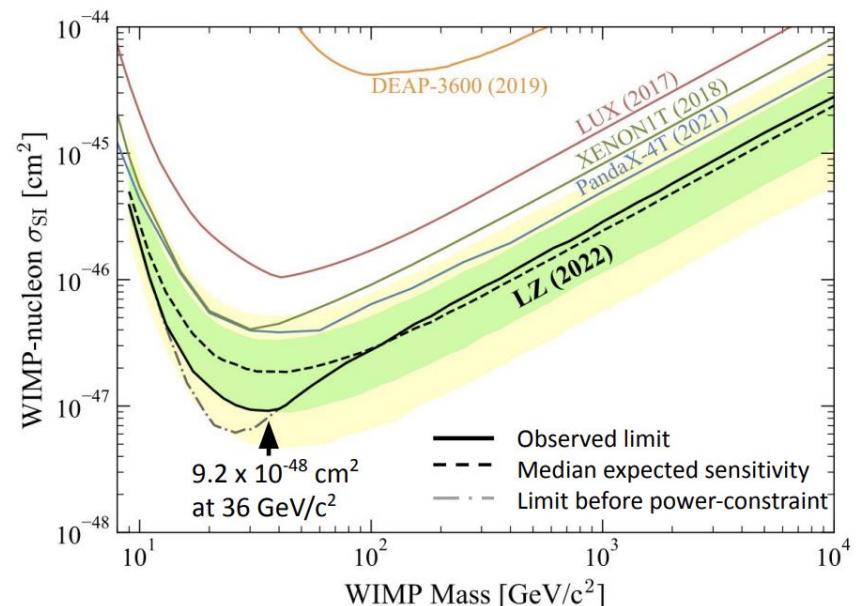
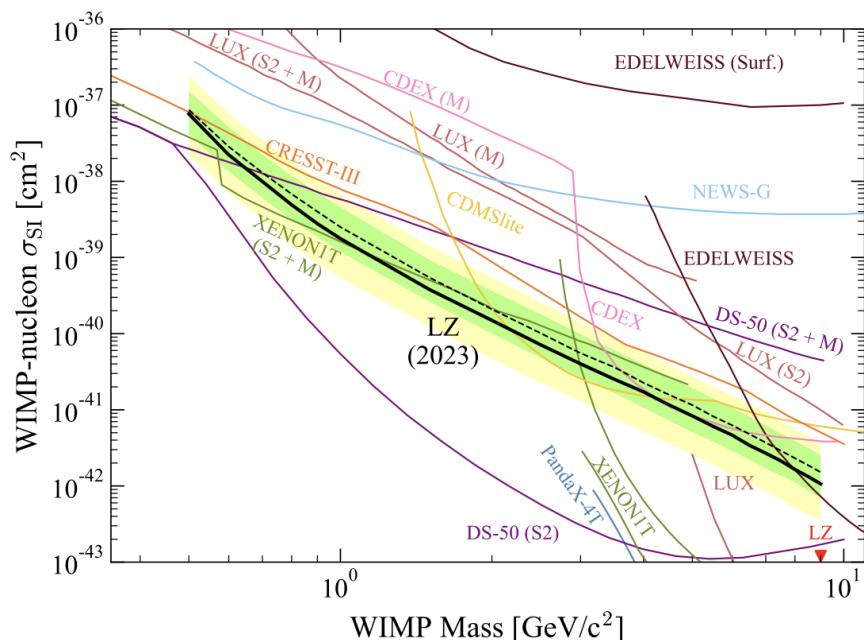
Dataset

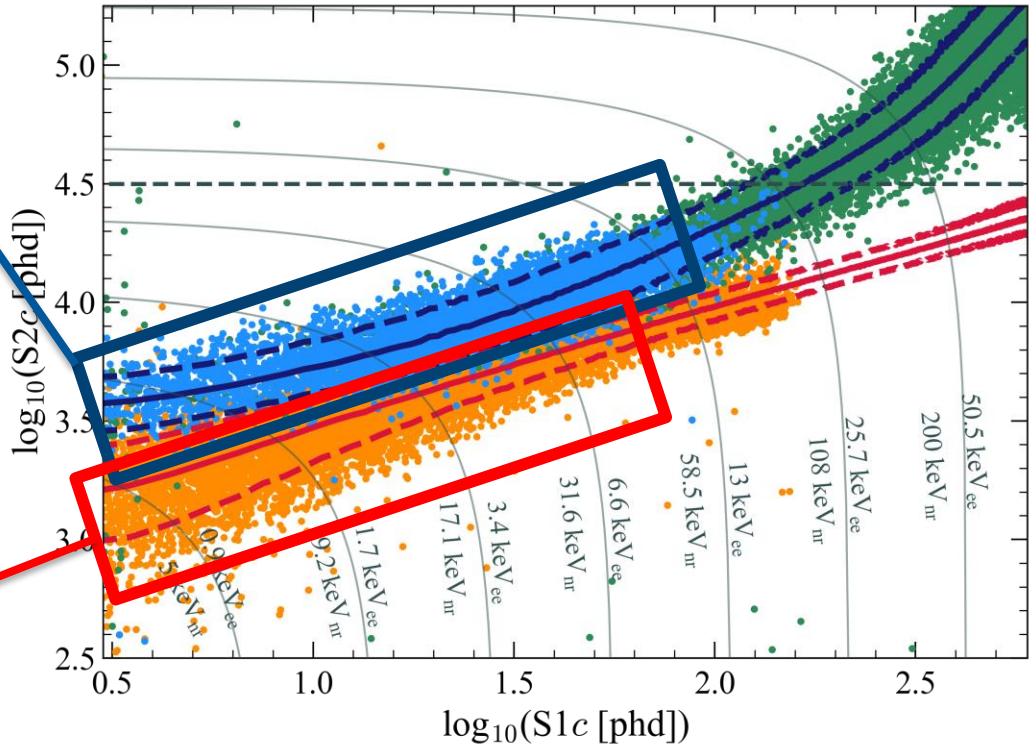
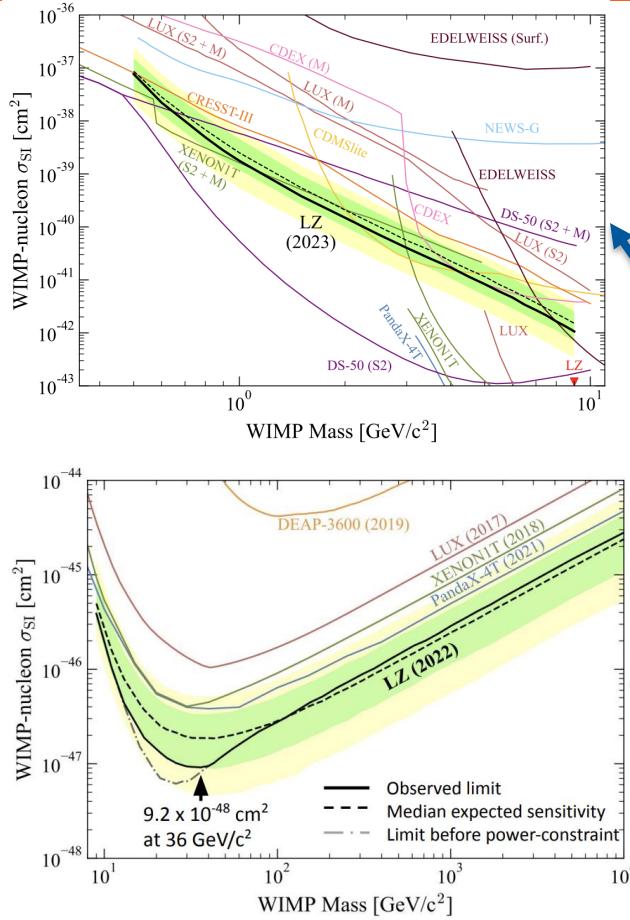
- Data taken Dec.2022-May.2023
- 60 ± 1 live days
- 5.5 ± 0.2 tonne fiducial volume
- Photon collection efficiency: $g1 = 0.114 \pm 0.2$ phd/photon
- Charge gain: $g2 = 47.1 \pm 1.1$ phd/electron

Analysis

- ER Calibrations (Tritium)
- NR Calibrations (Deuterium-Deuterium)
- Backgrounds well modelled in energy region (see [Phys. Rev. D 108, 012010](#))
- Unbinned profile likelihood in $\log_{10}(S2_c) - S1_c$
- Paper: [Phys. Rev. Lett. 131, 041002](#)









- Spin Independent and Spin Dependent interactions rely on the assumption of a zero-momentum transfer. But what if there is some momentum dependency?
- Use an EFT where we treat the WIMP-nucleon elastic scattering as a four-field interaction

$$\mathcal{L}_{int} = \mathcal{O} \chi^+ \chi^- N^+ N^-$$

- 4 Galilean, Hermitian invariants quantities which describe the interaction

$$i\vec{q}, \vec{S}_\chi, \vec{S}_N, \vec{v}^\perp \equiv \vec{v} + \frac{\vec{q}}{2\mu_N}$$

- We are then left with 15 operators which contribute to the interaction Lagrangian

$$\begin{aligned} \mathcal{L}_{int} &= \sum_i c_i \mathcal{O}_i \\ &= c_1 + ic_3 \vec{S}_N \cdot (\vec{q} \times \vec{v}^\perp) + c_4 \vec{S}_\chi \cdot \vec{S}_N \\ &\quad + ic_5 \vec{S}_\chi \cdot (\vec{q} \times \vec{v}^\perp) + c_6 (\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{q}) \\ &\quad + c_7 \vec{S}_N \cdot \vec{v}^\perp + c_8 \vec{S}_\chi \cdot \vec{v}^\perp + ic_9 \vec{S}_\chi \cdot (\vec{S}_N \times \vec{q}) \\ &\quad + c_{10} \vec{S}_N \cdot \vec{q} + ic_{11} \vec{S}_\chi \cdot \vec{q} + c_{12} \vec{S}_\chi \cdot (\vec{S}_N \times \vec{v}^\perp) \\ &\quad + ic_{13} (\vec{S}_\chi \cdot \vec{v}^\perp)(\vec{S}_N \cdot \vec{q}) + ic_{14} (\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{v}^\perp) \\ &\quad + -c_{15} (\vec{S}_\chi \cdot \vec{q}) ((\vec{S}_N \times \vec{v}^\perp) \cdot \vec{q}) \end{aligned}$$

- Differential Recoil Rate in this case

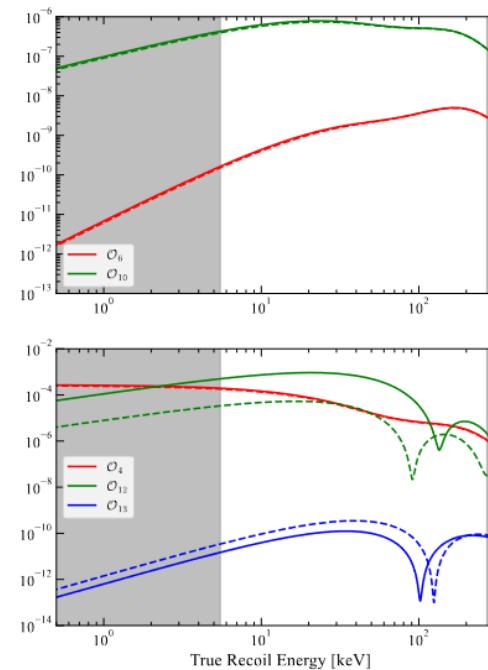
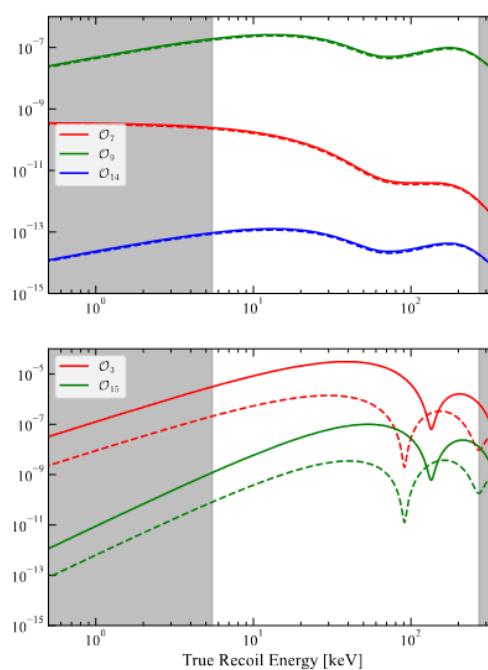
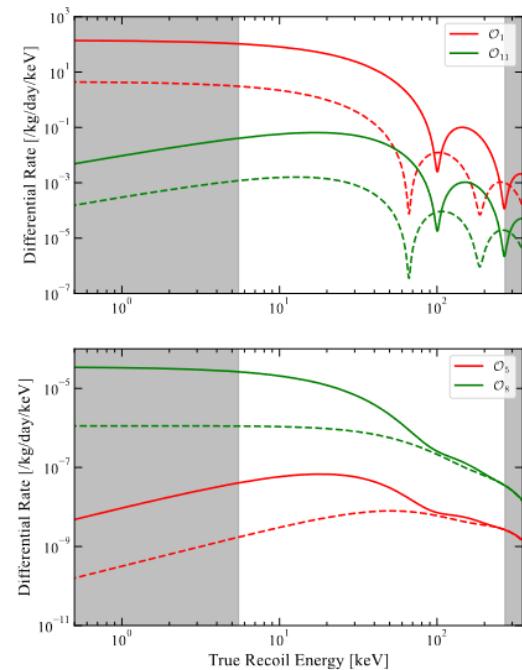
$$\frac{dR}{dE_R} = \frac{\rho_\chi}{32 \pi m_\chi^3 m_N^2} \int_{v>v_{min}}^\infty \frac{f(\vec{v})}{v} \sum_{i,j=1}^{15} \sum_{a,b=0,1} c_j^a c_i^b F_{i,j}^{a,b} d^3 v$$

[Phys. Rev. C 89, 065501 \(2014\)](#)



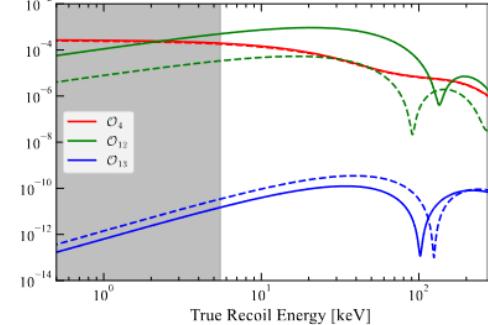
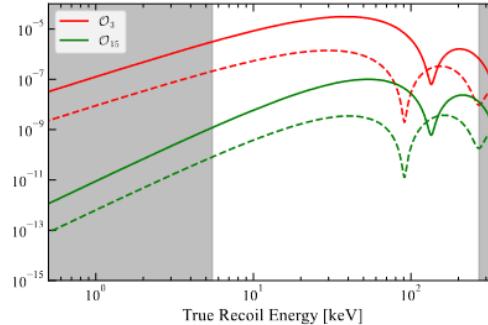
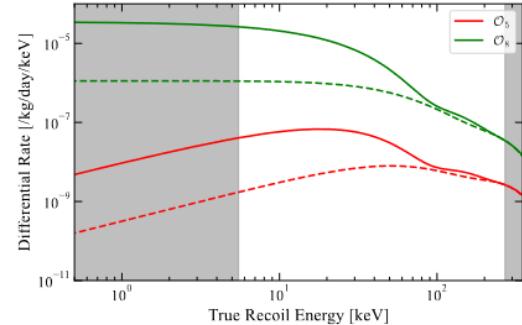
Evaluate the scattering amplitude assuming a single operator

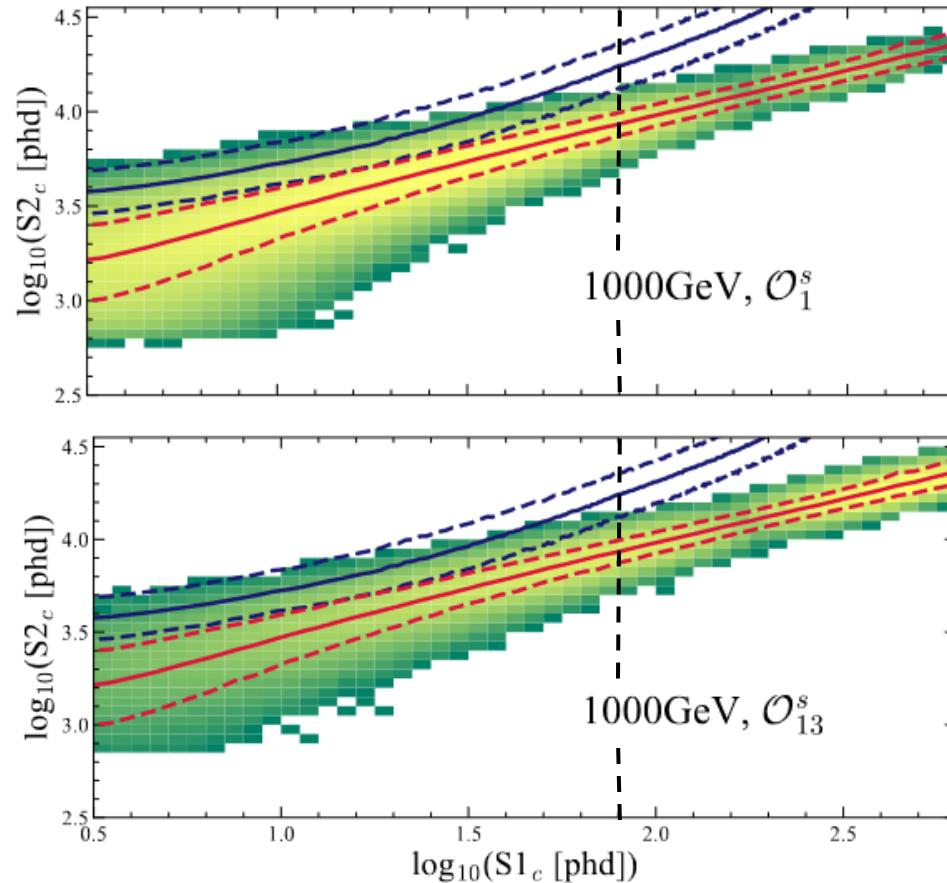
$$\frac{dR}{dE_R} \rightarrow \frac{\rho_\chi c_i^2}{32 \pi m_\chi^3 m_N^2} \int_{v > v_{min}}^\infty \frac{f(\vec{v})}{v} F_{i,i} d^3v$$



$m_\chi = 1000 \text{ GeV}/c^2$
Solid lines: isoscalar
Dashed lines: isovector

Shaded region:
Energy where the efficiency for the LZ
SR1 WIMP-Search data
is < 50%
After all cuts and ROI
selection







Can create an effective Lagrangian from combinations of various operators that themselves code in Galilean invariant interactions

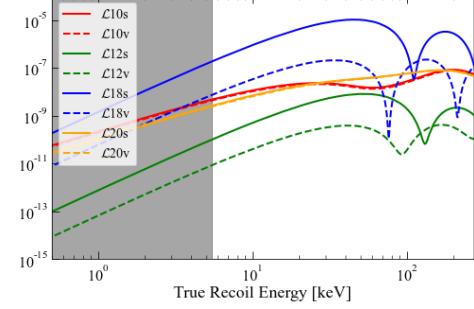
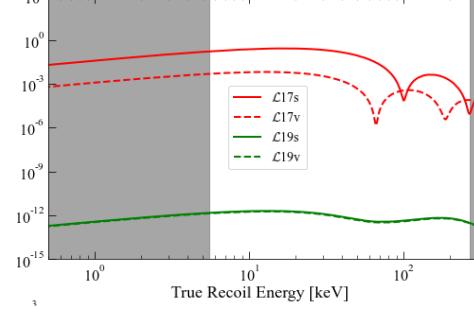
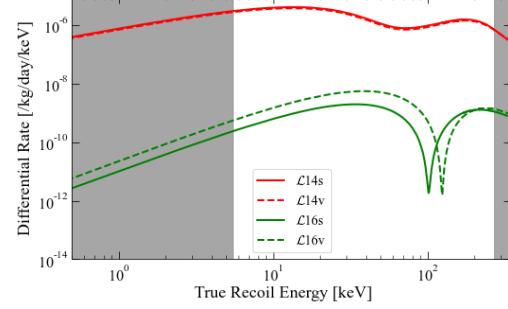
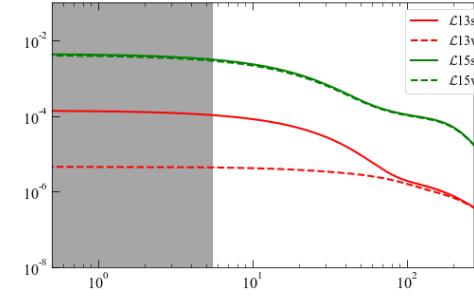
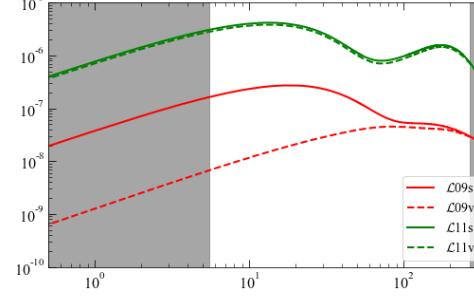
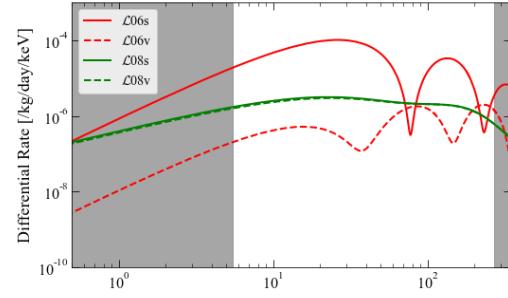
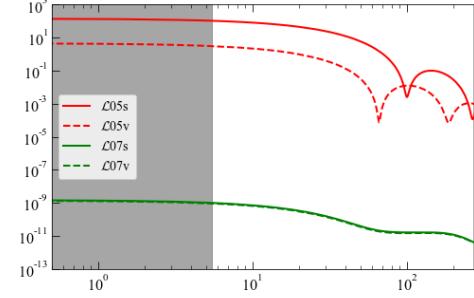
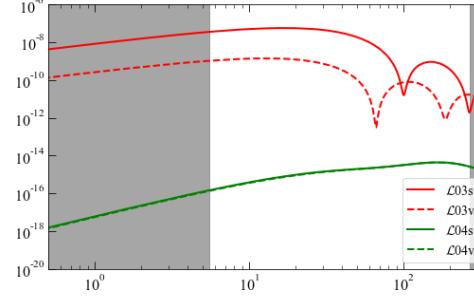
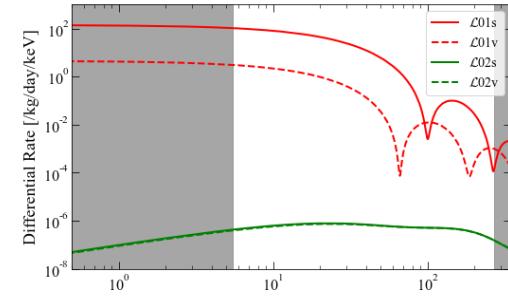
$$\mathcal{L}_{int} = \sum_{N=n,p} \sum_i d_i^{(N)} \mathcal{O}_i \bar{\chi} \chi \bar{N} N$$

d = operator construct coupling

WIMP magnetic moment: $\mathcal{L}_{int}^9 \rightarrow -\frac{\vec{q}^2}{2m_\chi m_M} \mathcal{O}_1 + \frac{2m_N}{m_M} \mathcal{O}_5 - \frac{2m_N}{m_M} \left(\frac{\vec{q}^2}{m_M} \mathcal{O}_4 - \mathcal{O}_6 \right)$

WIMP electric dipole moment: $\mathcal{L}_{int}^{17} \rightarrow \frac{2m_N}{m_M} \mathcal{O}_{11}$

[Phys. Rev. C 89, 065501 \(2014\)](#)



$m_\chi = 1000 \text{ GeV}/c^2$

Solid lines: isoscalar

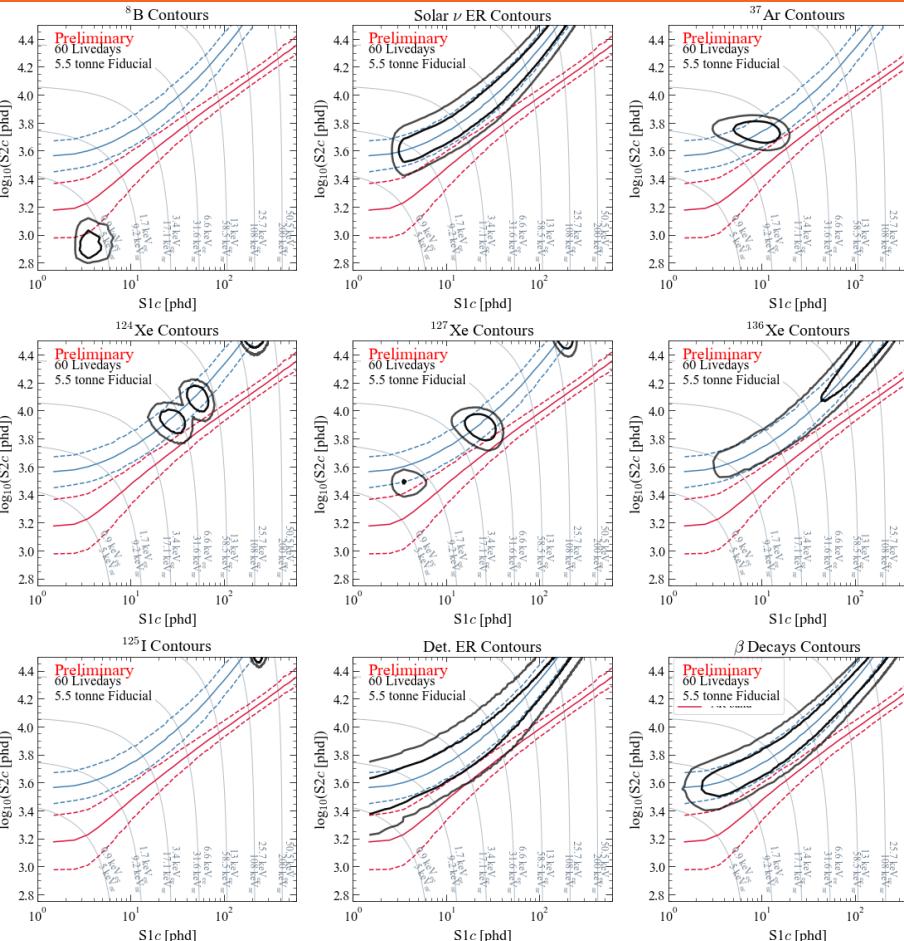
Dashed lines: isovector

Shaded region:

Energy where the efficiency for the LZ SR1 WIMP-Search data is <50%
After all cuts and ROI selection

All backgrounds relevant to SI WIMP search
Additional backgrounds

- Additional Xe decays
 - (^{124}Xe , ^{133}Xe , ^{131m}Xe , ^{127}Xe)
- Additional ER
- Topologies of more complex interactions
(eg multiple scatter signal ionization events)

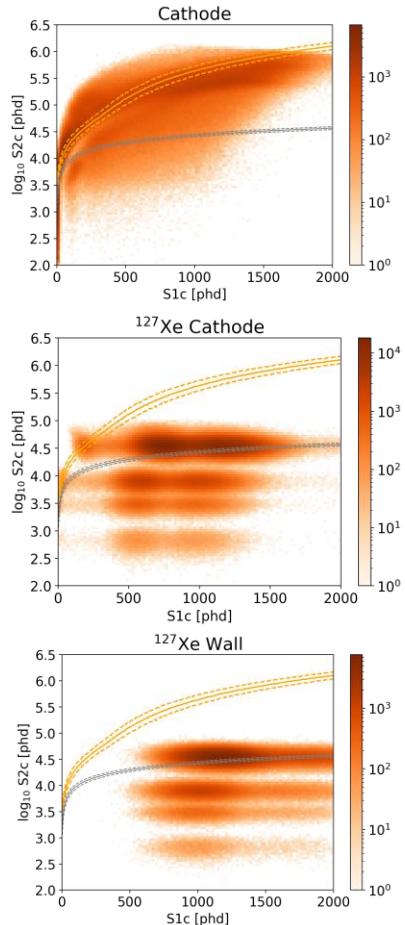
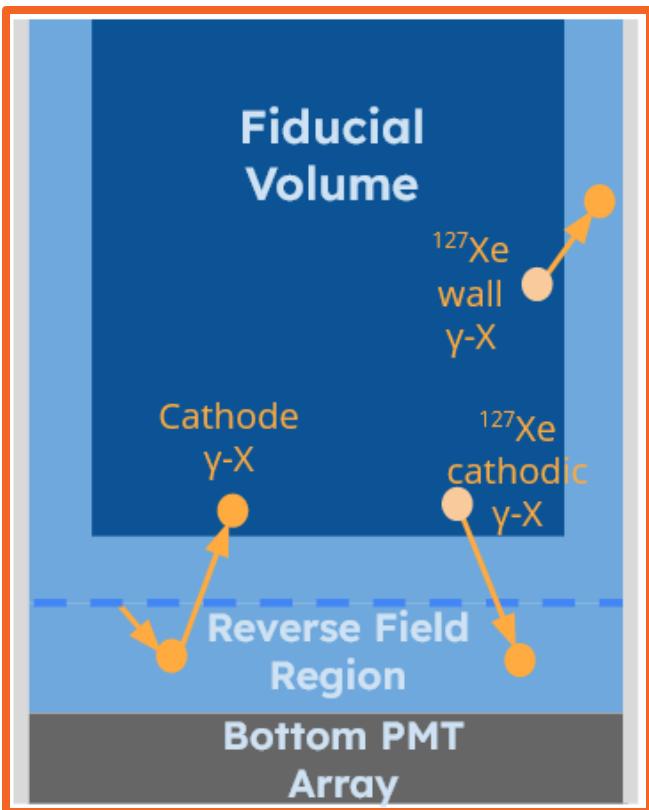


A γ -X event is a multiple-scattering γ where at least one vertex is in a region of incomplete charge collection:

- Reverse field region
- Near the TPC walls

Sources include:

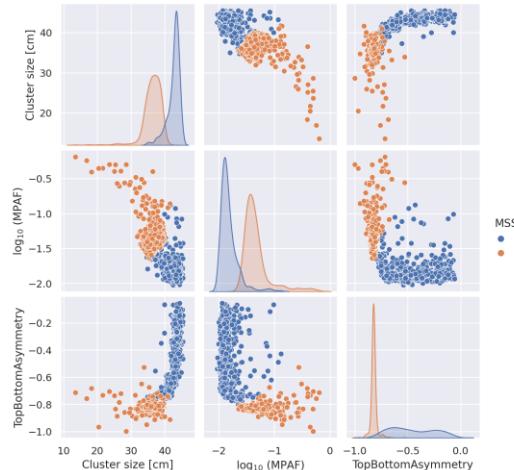
- ^{238}U , ^{232}Th , ^{60}Co , ^{40}K from cathode
- ^{127}Xe near detector edges





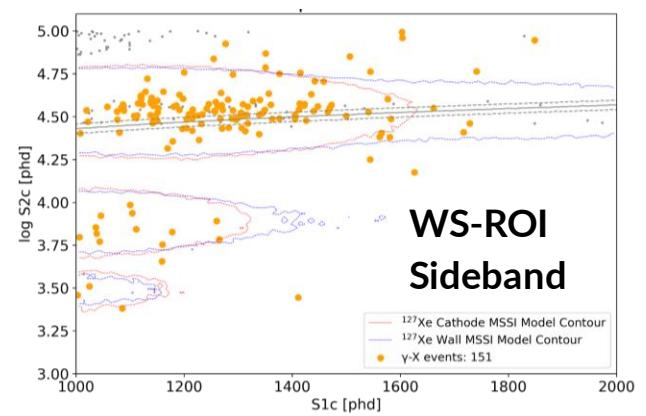
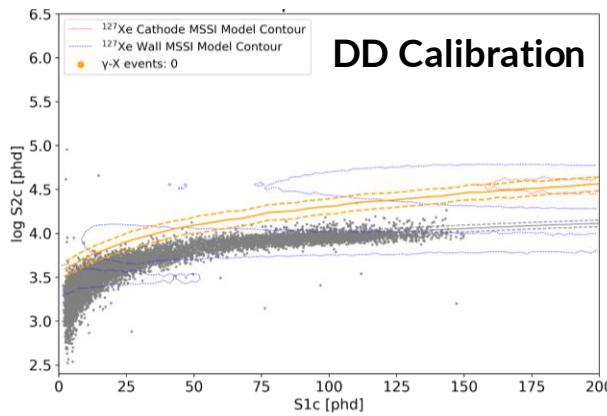
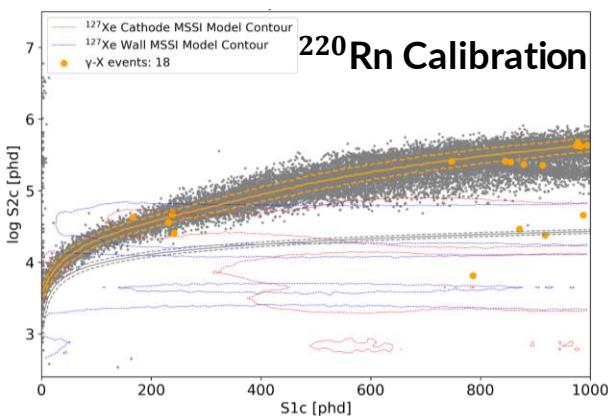
Boosted Decision Tree cut trained on high stats simulations and calibration data

$\downarrow P_e \rightarrow T$	SS	γ -X
SS	99.997 ± 0.005	0.4 ± 1.2
γ -X	0.003 ± 0.005	99.6 ± 1.2



Quantities used in the classification:

- Cluster size (size of the S1 splash on the bottom PMT array)
- Max Peak Area Fraction
- Top Bottom Asymmetry
- S1c
- \log_{10} S2c
- Radius
- Drift Time



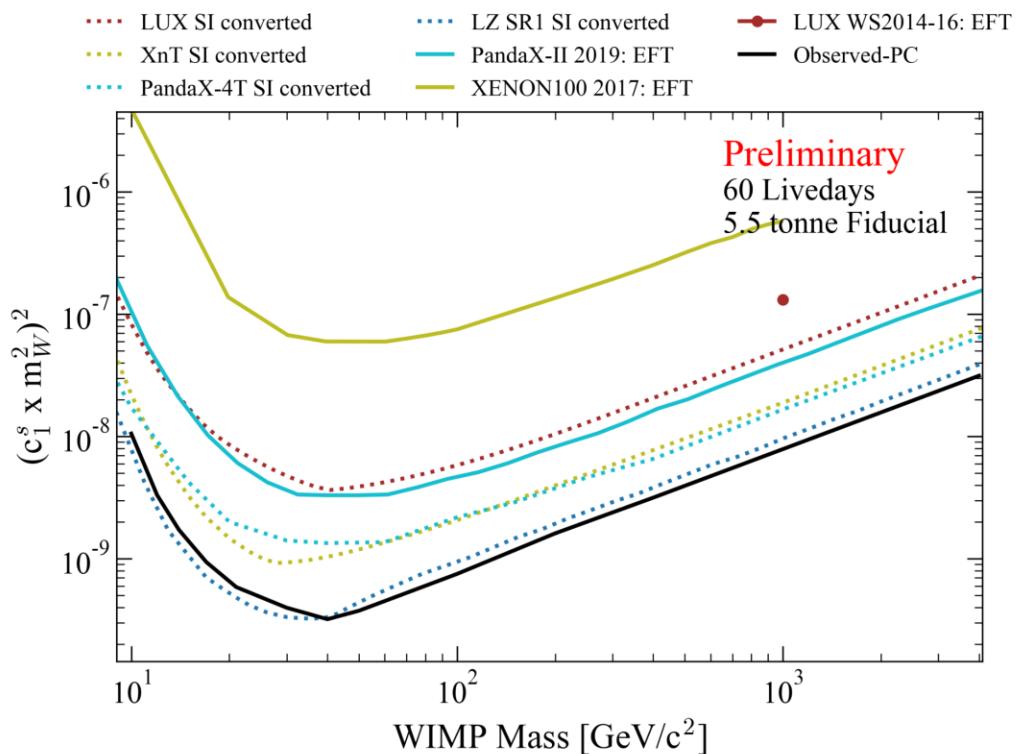


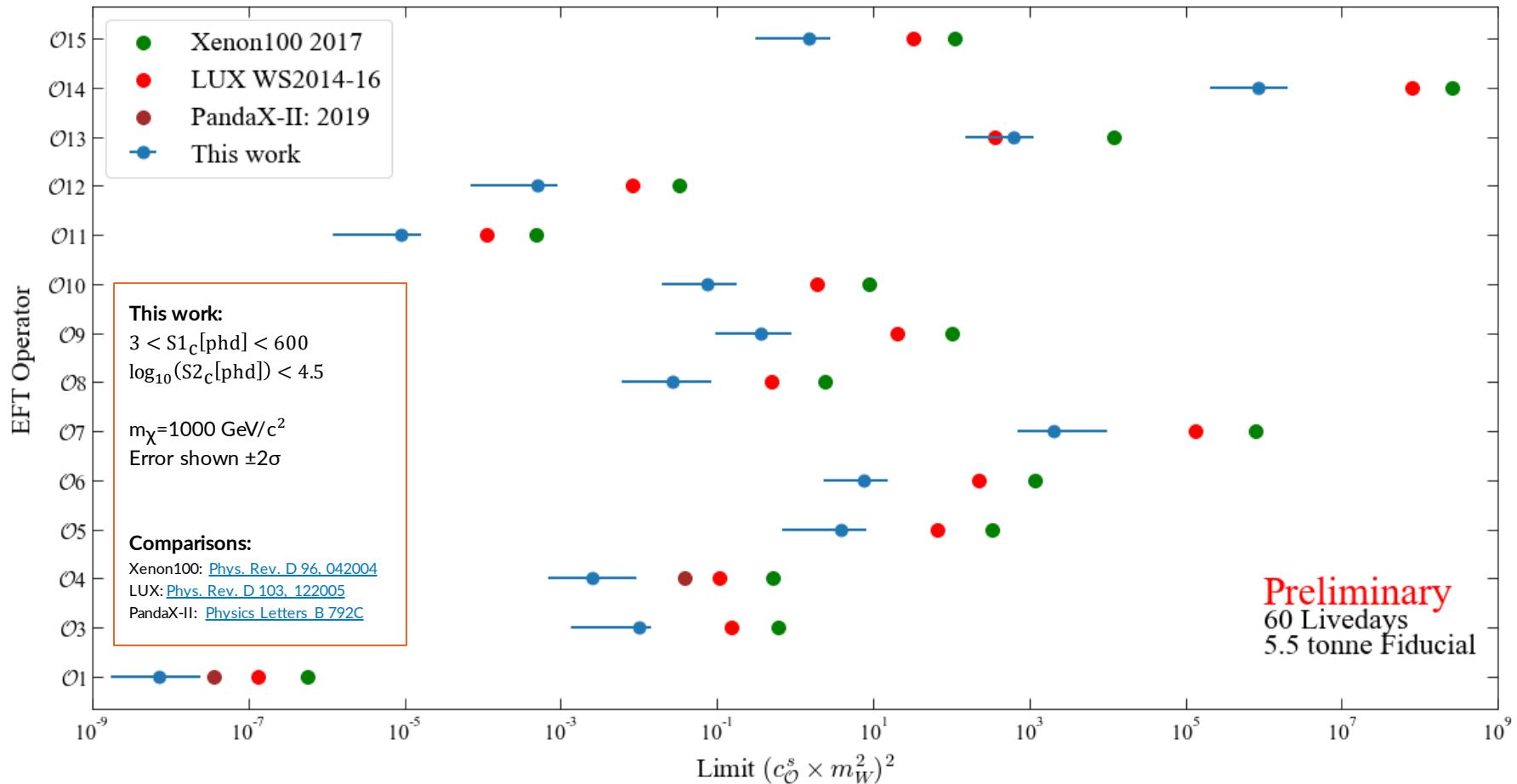
NREFT analyses often have differences in the choices of normalisations

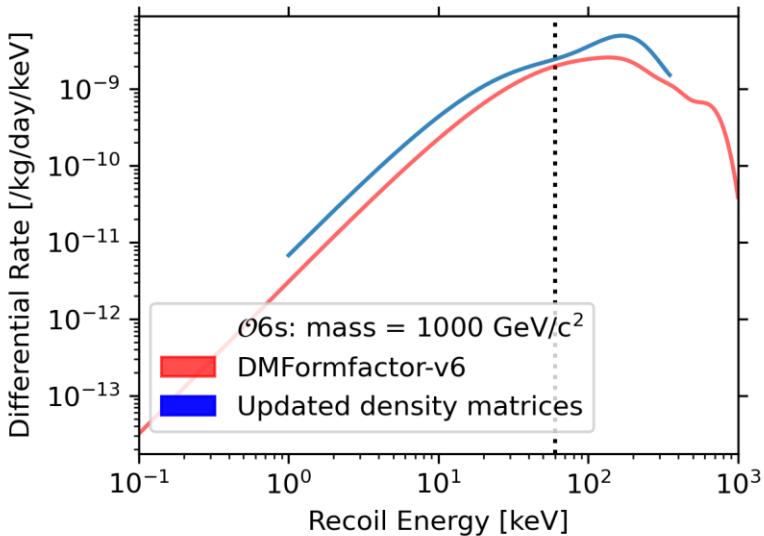
- Representation of isospin
- Dimensionality of the presented result

[Phys. Rev. D 96, 042004](#)

Experiment	Basis	Limit Type	Conversion in plot
Xenon100: 2017 EFT	$c_0 = \frac{1}{2}(c_p + c_n)$ $c_1 = \frac{1}{2}(c_p - c_n)$	$(c_1^s \times m_w^2)^2$	None
LUX: WS2014-16 EFT	$c_0 = (c_p + c_n)$ $c_1 = (c_p - c_n)$	$(c_1^s \times m_w^2)^2$	$\frac{1}{4}$
Physics Letters B 792C	PandaX-II: SD EFT	$c_0 = \frac{1}{2}(c_p + c_n)$ $c_1 = \frac{1}{2}(c_p - c_n)$	$d_5^s \left[\frac{1}{m_w^2} \right]$ $(d_5^s)^2$
Paper in progress	LZ EFT (This analysis)	$c_0 = \frac{1}{2}(c_p + c_n)$ $c_1 = \frac{1}{2}(c_p - c_n)$	$(c_1^s \times m_w^2)^2$
Phys. Rev. C 89, 065501	NRET Theory paper	$c_0 = \frac{1}{2}(c_p + c_n)$ $c_1 = \frac{1}{2}(c_p - c_n)$	N/A N/A
Phys. Rev. Lett. 118, 021303	LUX: Combined 2017 SI	N/A	σ_{SI}^N
Phys. Rev. Lett. 127, 261802	PandaX-4T: 2021 SI	N/A	$\sigma_{SI}^N \frac{\pi \cdot m_w^4}{(\frac{(hc)}{GeV})^2 \mu_N^2}$
Phys. Rev. Lett. 131, 041002	LZ: 2023 SI	N/A	$\sigma_{SI}^N \frac{\pi \cdot m_w^4}{(\frac{(hc)}{GeV})^2 \mu_N^2}$
Phys. Rev. Lett. 131, 041003	XENONnT: 2023 SI	N/A	$\sigma_{SI}^N \frac{\pi \cdot m_w^4}{(\frac{(hc)}{GeV})^2 \mu_N^2}$

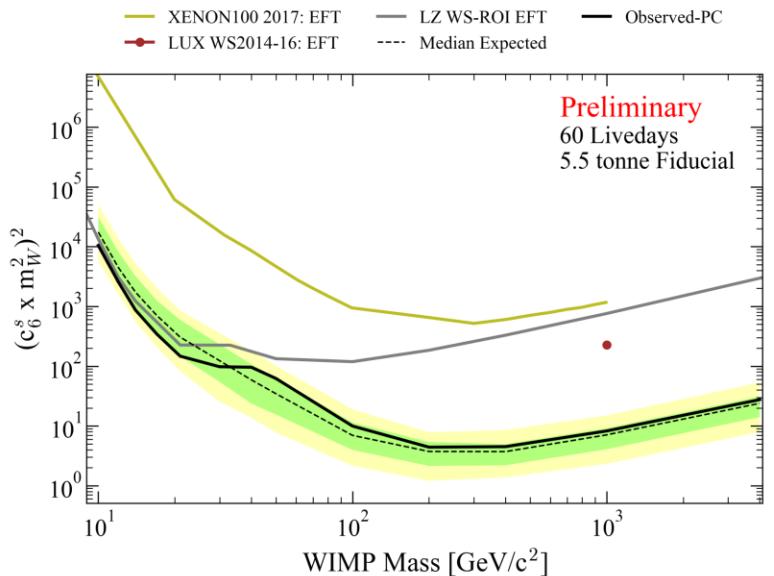






Dashed line indicates the energy at falling edge of the efficiency is 50% with LZ SR1 WIMP-Search ROI:

- $3 < S1_C[\text{phd}] < 80$



Grey:

- $3 < S1_C[\text{phd}] < 80$
 - DMFormFactor-v6 signal
- Black + **Brazilian band** (this work):
- $3 < S1_C[\text{phd}] < 600, \log_{10}(S2_C[\text{phd}]) < 4.5$
 - Updated density matrices for signals

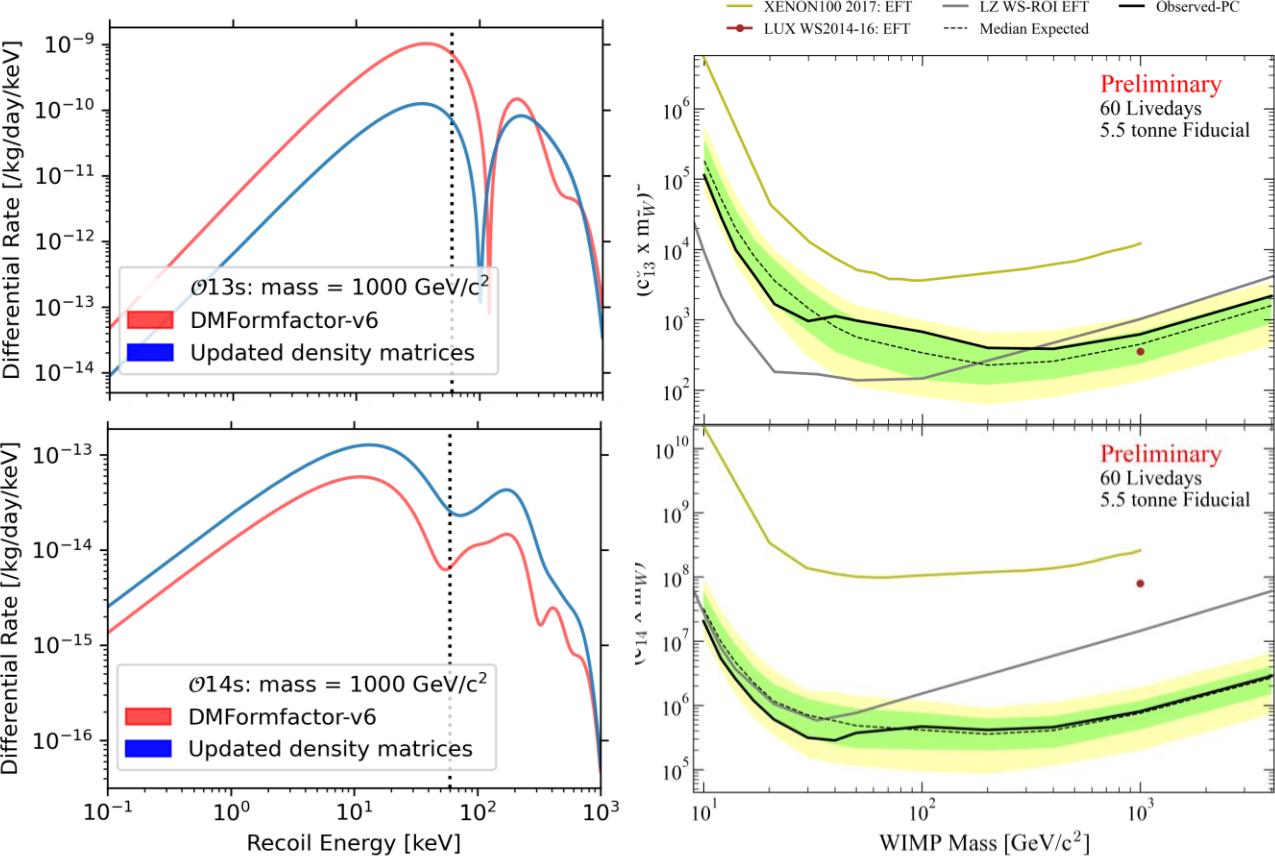


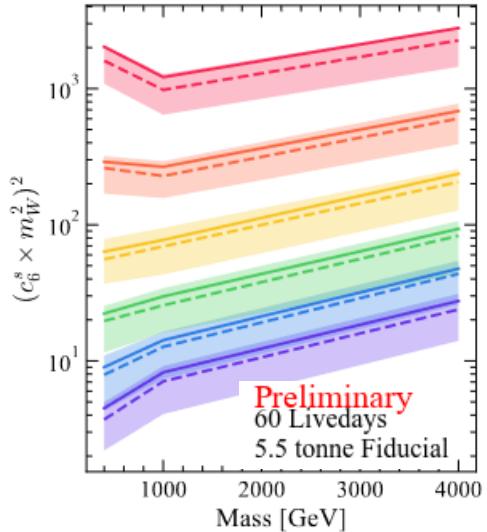
Differential recoils used updated GCN5082 ground state to ground state one-body density matrices

These have significant impacts on some operators

\mathcal{O}_{13} differential rate has decreased significantly

\mathcal{O}_{14} has the opposite behaviour





Inelastic Operators

This work:

- $3 < S1_C[\text{phd}] < 600$
- Updated density matrices for signals

$$\delta_m \equiv m_{\chi,out} - m_{\chi,in}$$

$$\delta_m + \vec{v} \cdot \vec{q} + \frac{|\vec{q}|^2}{2\mu_N} = 0$$

$$\vec{v}_{inel}^\perp \equiv \vec{v} + \frac{\vec{q}}{2\mu_N} + \frac{\delta_m}{|\vec{q}|^2} \vec{q}$$

Elastic Lagrangians

This work:

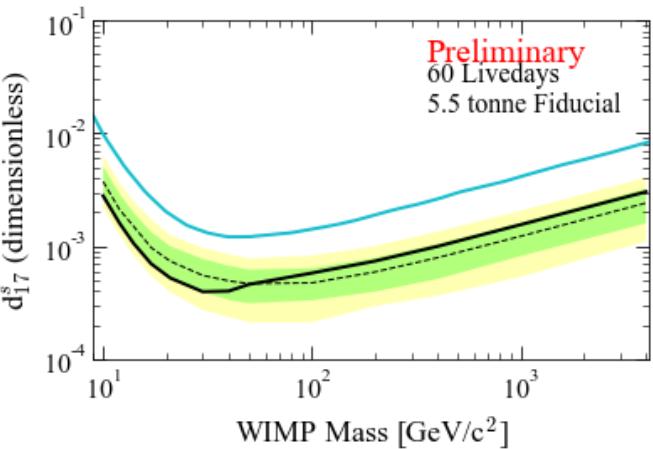
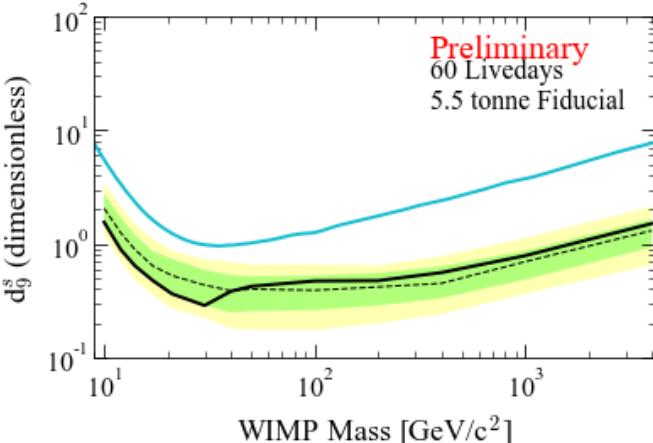
- $3 < S1_C[\text{phd}] < 600$
- Updated density matrices for signals

Blue:

- PandaX-II: [Physics Letters B 792C](#)

$$\mathcal{L}_{int}^9 \rightarrow -\frac{\vec{q}^2}{2m_\chi m_M} \mathcal{O}_1 + \frac{2m_N}{m_M} \mathcal{O}_5 - \frac{2m_N}{m_M} \left(\frac{\vec{q}^2}{m_M} \mathcal{O}_4 - \mathcal{O}_6 \right)$$

$$\mathcal{L}_{int}^{17} \rightarrow \frac{2m_N}{m_M} \mathcal{O}_{11}$$





- Initial science run of LZ has placed stringent limits on both SI and SD WIMP-nucleon interactions
- These models are relatively simplistic, and the inherent nature of interaction may be more complex
- Using an EFT allows us to describe all possible dark matter interactions with nucleons
- Extending energy region is beneficial for EFT analysis
- LZ has set promising EFT limits in SR1 with an extended energy ROI
- With the energy region understood in LZ, we can test a lot of DM parameter space, eg:
 - 2HDM+a: [Phys. D. M. 27, 100351 \(2020\)](#)
 - DM -photon interactions: [Nature 618, 47 \(2023\)](#)

LZ Publications:

- Detector paper: [Nucl. Instrum. Meth. A 953, 163047](#)
- SR1 papers:
 - WIMP-nucleon result: [Phys. Rev. Lett. 131, 041002](#)
 - Backgrounds: [Phys. Rev. D 108, 012010](#)
 - LowE ER: [arXiv:2307.15753](#)
 - this work: paper in progress



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Questions?



Science and
Technology
Facilities Council

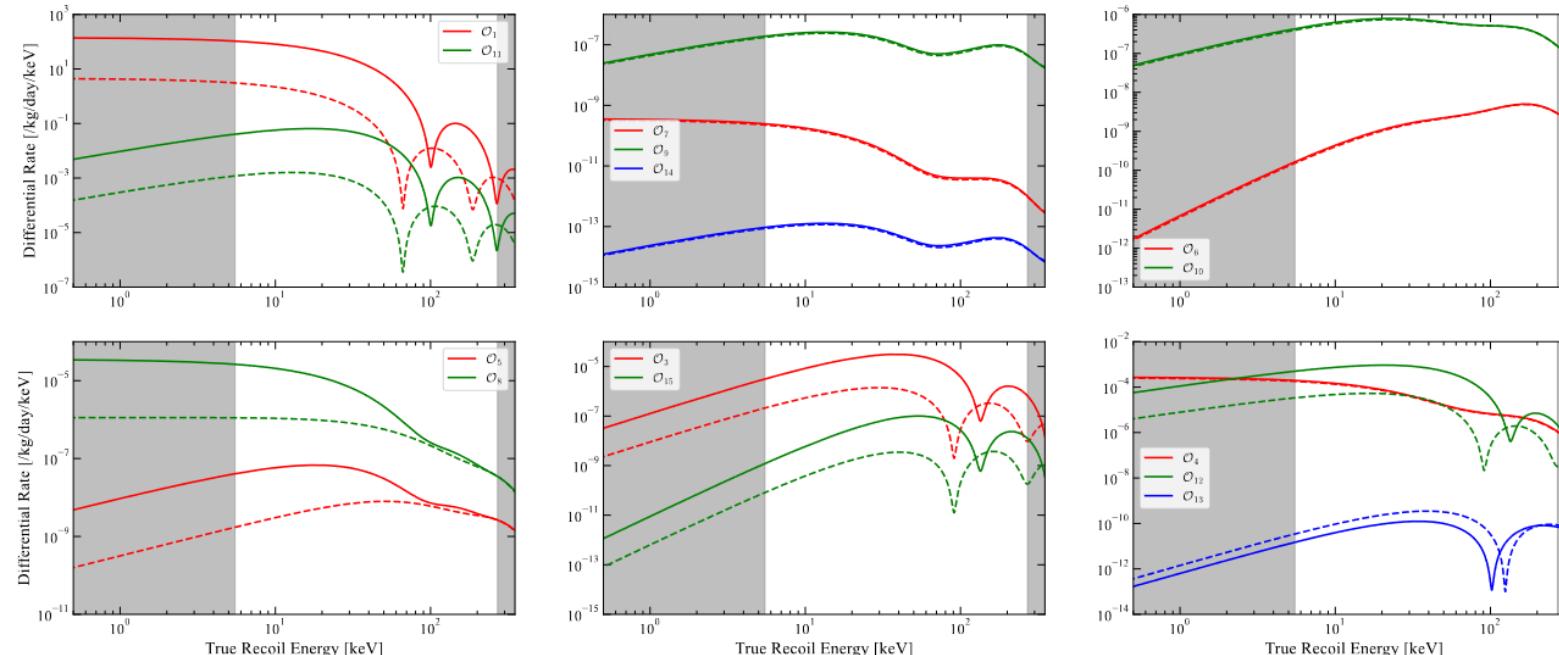


FCT

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MINISTÉRIO DA EDUCAÇÃO E CIÉNCIA

ibS Institute for
Basic Science





Code used:

[DMFormFactor-v6](#): Used for updating the nuclear response function

[WimPyDD](#): Used to generate final recoils with updated response functions

Using updated GCN5082 ground state to ground state one-body density matrices ([supplied](#) by W. Haxton, generated using [BIGSTICK](#))

Evaluate the scattering amplitude assuming a single operator

$$\frac{dR}{dE_R} \rightarrow \frac{\rho_\chi c_i^2}{32 \pi m_\chi^3 m_N^2} \int_{v>v_{min}}^\infty \frac{f(\vec{v})}{v} F_{i,i} d^3 v$$



Galilean-invariant to quadratic order in momentum transfer

$$i\vec{q}, \vec{S}_\chi, \vec{S}_N, \vec{v}^\perp \equiv \vec{v} + \frac{\vec{q}}{2\mu_N}$$

Spin 1 or less particles

Theory translated to coefficients of an effective operator

These operators can be equated to generic nuclear responses

Possible to reduce covariant interaction Lagrangians to combinations of the operators

$$\frac{dR}{dE_R} = \frac{\rho_\chi}{32 \pi m_\chi^3 m_N^2} \int_{v>v_{min}}^\infty \frac{f(\vec{v})}{v} \sum_{i,j=1}^{15} \sum_{a,b=0,1} c_j^a c_i^b F_{i,j}^{a,b} d^3 v$$

Interactions are linear combinations of 6 independent nuclear responses

M Spin Independent

Σ' Spin Dependent (transverse)

Σ'' Spin Dependent (longitudinal)

Δ Angular Momentum

Φ'' Spin Orbit

$\tilde{\phi}'$ Tensor spin orbit

Operators give us a sense of the true sensitivity of our detector to different DM-nucleon physics

[Phys. Rev. C 89, 065501 \(2014\)](#)